



## Research Article

# Patterned growth of $\beta$ -Ga<sub>2</sub>O<sub>3</sub> thin films for solar-blind deep-ultraviolet photodetectors array and optical imaging application

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## ABSTRACT

Solar-blind deep-ultraviolet (DUV) photodetectors based on Ga<sub>2</sub>O<sub>3</sub> have attracted great attention due to their potential applications for many military and civil purposes. However, the development of device integration for optoelectronic system applications remains a huge challenge. Herein, we report a facile method for patterned-growth of high-quality  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> thin films, which are assembled into a photodetectors array comprising 8 × 8 device units. A representative detector exhibits outstanding photoresponse performance, in terms of an ultra-low dark current of ~0.62 pA, a large  $I_{\text{light}}/I_{\text{dark}}$  ratio exceeding 10<sup>4</sup>, a high responsivity of ~0.72 A W<sup>-1</sup> and a decent specific detectivity of ~4.18 × 10<sup>11</sup> Jones, upon 265 nm DUV illumination. What is more, the DUV/visible (250/400 nm) rejection ratio is as high as 10<sup>3</sup> with a sharp response cut-off wavelength at ~280 nm. Further optoelectronic analysis reveals that the photodetectors array has good uniformity and repeatability, endowing it the capability to serve as a reliable DUV light image sensor with a decent spatial resolution. These results suggest that the proposed technique offers an effective avenue for patterned growth of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> thin films for multifunctional DUV optoelectronic applications.

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## Introduction

Solar-blind deep-ultraviolet (DUV) spectrum regime refers to the solar radiation with wavelength of 200–280 nm, which does not exist at the surface of the Earth because of the strong absorption of ozone and atmosphere layers [1,2]. Photodetectors operating in this waveband region with a response cut-off wavelength around or below 280 nm are named as solar-blind photodetectors, which have pivotal applications in a variety of military and civil areas including missile guidance systems, intersatellite communication, flame detection, biochemical detection and remote control, etc [3,4]. Holding apparent advantages such as large bandgaps significantly exceeding the 3.4 eV of GaN, intrinsic solar-blindness and high radiation hardness, ultra-wide bandgap (UWBG) semiconductors including Al<sub>x</sub>Ga<sub>1-x</sub>N, BN, Zn<sub>x</sub>Mg<sub>1-x</sub>O, gallium oxide (Ga<sub>2</sub>O<sub>3</sub>) and diamond are promising candidates for solar-blind DUV photodetection, and have also been extensively explored in recent years [2,5,6].

Among them, monoclinic Ga<sub>2</sub>O<sub>3</sub> ( $\beta$ -Ga<sub>2</sub>O<sub>3</sub>), which is the most stable form of Ga<sub>2</sub>O<sub>3</sub> under normal conditions, has a direct bandgap ranging from 4.4–5.1 eV, corresponding to an absorption cut-off wavelength in the range of 240–280 nm [7]. Such an optical property enables photodetection over the full region of solar-blind DUV. Moreover, this kind of material exhibits high thermal and chemical stabilities [8]. The above characteristics make  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> extremely attractive for developing solar-blind DUV photodetectors, in particular, with the purpose for use in some high-temperature or harsh environments [9,10]. So far, a large number of high-performance DUV photodetectors composed of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> in bulk, thin film or nanostructure forms have been reported [11–17]. Different device geometries including photoconductors [13], metal-semiconductor-metal (MSM) photodetectors [18], heterojunction/Schottky barrier photodiodes [19–21], and avalanche photodiodes [22,23] have also been involved. Among these, MSM photodetectors based on  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> thin films are of particular interest due to a range of features such as mature material preparation techniques, simple device fabrication, low dark current, the potential to attain high photogains with large responsivities, the good homogeneity of the material for easy integration with readout circuitry, etc [24].

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By far,  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> thin film-based MSM photodetectors in the form of individual device unit have been widely studied [2,11]. Most recently, several research groups have presented two-dimensional (2D) or even three-dimensional (3D) photodetectors array based on monolithic  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> thin films that exhibit outstanding DUV photoresponse with good device uniformity. [24–26] In addition, preliminary device integration technique towards multifunctional optoelectronic applications such as DUV image sensing, multipoint light spatial distribution recognition and real-time light trajectory has also been successfully developed. From the perspective of device fabrication, these applications often involve the patterning of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> thin films by a controlled etching process, which is essentially required to achieve intra-device isolation and depress crosstalk between adjacent devices [7]. Thin films of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> are normally grown on sapphire substrates via synthetic methods such as molecular beam epitaxy (MBE), metalorganic chemical vapor deposition (MOCVD), magnetron sputtering and so on [27,28]. To attain a good crystalline quality, high temperature is usually involved in the preparation process, which precludes the possibility of pre-patterning the material assisted by a photolithography process [29–31]. On the other hand, etching is also problematic for  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> material due to the difficulties such as isotropic etching and surface dangling bonds-dependent etching properties in wet etching, as well as low etching rate, poor surface morphology and introduction of ion-induced damage in dry etching [7,32,33].

Herein, we present a technique combing procedures of standard ultraviolet (UV) photolithography, metal deposition and thermal-assisted conversion, for patterned growth of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> thin films with high crystalline quality. Via this method, MSM photodetectors array composed of 8 × 8 device units are realized, which displays not only a good photoresponse to DUV irradiation but also a narrow unit-to-unit variation. The important performance parameters including dark current,  $I_{\text{light}}/I_{\text{dark}}$  ratio, responsivity and specific detectivity attain  $\sim 0.62$  pA,  $\sim 6.13 \times 10^4$ ,  $\sim 0.72$  A W<sup>-1</sup> and  $\sim 4.18 \times 10^{11}$  Jones, respectively, at 265 nm DUV illumination. In addition, the detectors also display a steep response cut-off wavelength at around 280 nm with the DUV/visible (250/400 nm) rejection ratio exceeding 10<sup>3</sup>. Furthermore, the present patterned photodetectors array also shows the capability to serve as a DUV light image sensor with satisfying resolution. It is believed that our study provides a facile, reliable and low-cost pathway for patterned growth of high-quality  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> thin films towards device integration for multifunctional DUV optoelectronic application purposes.

## Experimental

### Material preparation and characterization

The patterned  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> thin films were prepared by using a method containing standard ultraviolet (UV) photolithography, metal deposition and thermal-assisted conversion processes. A SiO<sub>2</sub>/Si substrate was ultrasonically cleaned sequentially in acetone, ethanol and deionized water for 10 min, respectively, and dried in a flow of N<sub>2</sub> gas prior to use. The substrate was then spin-coated with a solution of positive photoresist, and a standard UV photolithography process was performed, which defined the photoresist layer with designed patterns. Afterwards, Ga layer with a thickness of  $\sim 22$  nm was deposited onto the substrate in an ultrahigh vacuum chamber via thermal evaporation of metallic Ga (Aldrich, 99.99%). After a lift-off process, patterned Ga thin films was formed on the substrate. Subsequently, the substrate was transferred into a horizontal tube furnace. After evacuation, a constant flow of O<sub>2</sub> gas (50 sccm) was introduced through the tube, and the tube pressure was maintained at 50 mTorr. The furnace was then heated to 1100 °C and kept at this temperature for 1 h

to conduct the thermal-assisted conversion. After reaction, the furnace was cooled down naturally, and the patterned  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> thin films were obtained on the substrate.

The morphology and composition of the as-prepared  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> thin films were studied by a field-emission scanning electron microscope (FESEM, Quanta, FEG 250) equipped with an energy-dispersive spectrometer, and X-ray diffraction (XRD, X'Pert PRO MPD). The absorption spectrum was recorded on a CARY 5000 UV–vis (NIR) spectrophotometer. The X-ray photoelectron spectroscopy (XPS) measurement was performed by using a monochromatic Al K $\alpha$  source produced by the XPS system (Thermo ESCALAB250 X-ray photoelectron spectroscopy).

### Device fabrication and analysis

The  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> thin film-based photodetectors array was constructed following the procedures as schematically shown in Fig. 1.  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> thin films with the pattern of an 8 × 8 squares array was first obtained on a cleaned SiO<sub>2</sub>/Si substrate by using the method described above. The size of an individual square pattern is 300  $\mu$ m × 300  $\mu$ m. Then, Au electrodes (50 nm) array with a channel length of 45  $\mu$ m was deposited by another UV photolithography, electron beam deposition and lift-off processes. The electrical properties was characterized by using a semiconductor parameter analyzer (Keithley 4200-SCS) on a probe station. A UV light-emitting diode (LED) with the wavelength of 265 nm (Thorlabs M265L3) serving as the light source was used to study the photoresponse performance. For spectral response study, a lab-built optoelectronic system consisting of a light source (LE-SP-LS-XE) and monochromator (LE-SP-M300) was employed. Before use, the intensity of all light sources was carefully calibrated by using a power meter (Thorlabs GmbH, PM 100D). All measurements were carried out in air at room temperature.

## Results and discussion

In this work, the patterned  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> thin films were prepared through a thermal-assisted conversion method, as schematically shown in Fig. 1. The detailed procedures can be found in the “Experimental Section”. First, a standard UV photolithography was carried out to define a photoresist layer on a cleaned SiO<sub>2</sub>/Si substrate with predesigned patterns. Then, the patterned Ga thin films were obtained on the substrate through thermal evaporation, followed by a lift-off process. Finally,  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> thin films with patterns were formed by transferring the substrate into a tube furnace, where a thermal-assisted conversion process was conducted. To fabricate the  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> thin film-based photodetectors array, another procedure containing UV photolithography, Au deposition via electron-beam evaporation and lift-off was then performed.

The morphological and structural characterizations were carried out to study the quality of the as-grown patterned  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> thin films. Fig. 2(a)–(c) illustrates the FESEM images of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> thin films of periodic square array, round array and triangle array. Apparently, these films exhibited uniform and regular shapes with desired patterns, implying the successful patterned growth of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> thin films with the proposed technique. As observed from a magnified FESEM image (insert in Fig. 2(b)), the  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> thin films were virtually composed of closely compact crystalline grains with size of  $\sim 100$  nm, similar to those  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> thin films grown by plasma enhanced chemical vapor deposition (PECVD) technique [20]. In addition to the above patterns, our technique could also be applied to grow  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> thin films with other patterns. As shown in Fig. 2(d), stripe patterns with a gap width of 10  $\mu$ m were realized, and even with the stripe width of only 4  $\mu$ m, continuous and regular stripes of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> thin film with sharp edges could still

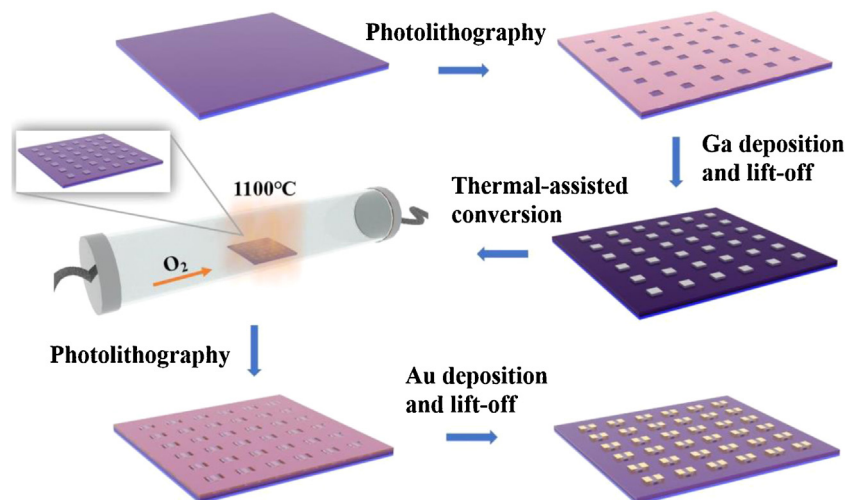


Fig. 1. Schematic illustration of the procedures for fabricating patterned  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> thin film-based photodetectors array.

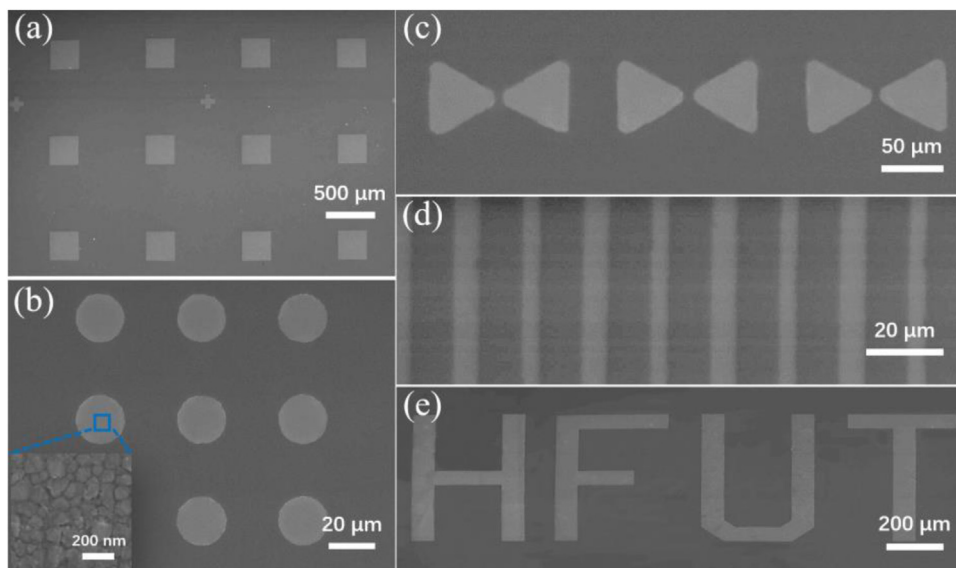


Fig. 2. FESEM images of the  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> thin films with the patterns of (a) periodic square array, (b) periodic round array, (c) periodic triangle array, (d) continuous 4 and 6  $\mu$ m stripes and (e) "HFUT" characters. Inset in (b) shows an enlarged FESEM image of the  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> thin film.

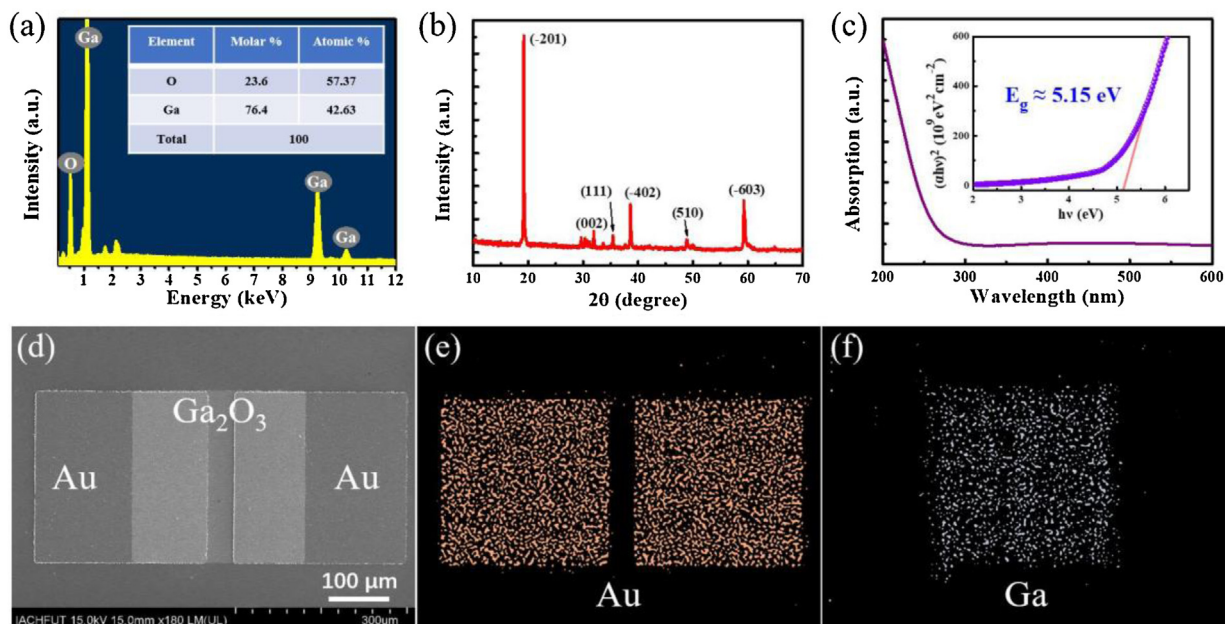
be obtained. Moreover, Fig. 2(e) displays a FESEM image of patterned  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> thin films arranged in a shape of "HFUT" character. Besides, it should be mentioned that un-patterned  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> thin films in a large scale could also be grown using our method without the process of UV photolithography (Fig. S1 in Supporting Information). These results suggested that the proposed technique held high maneuverability and was scalable for growing  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> thin films with more complicated patterns towards multiple applications.

Fig. 3(a) depicts the X-ray energy-dispersive spectrometer (EDS) profile of the as-prepared thin films, from which the atomic ratio of Ga:O was determined to be 42.63:57.37. This value was in good accordance with the result obtained from XPS analysis (Fig. S2 in Supporting Information), and was close to the stoichiometric ratio of Ga<sub>2</sub>O<sub>3</sub> as well. The XRD pattern plotted in Fig. 3(b) displays several evident peaks, which could be readily assigned to the  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> (JCPDS card No. 43-1012) [15]. Furthermore, from the optical absorption spectrum of the films (Fig. 3(c)), pronounced absorption at the DUV wavelength regime ranging from 200 to 280 nm along with a steep absorption cut-off edge situated at  $\sim$ 280 nm

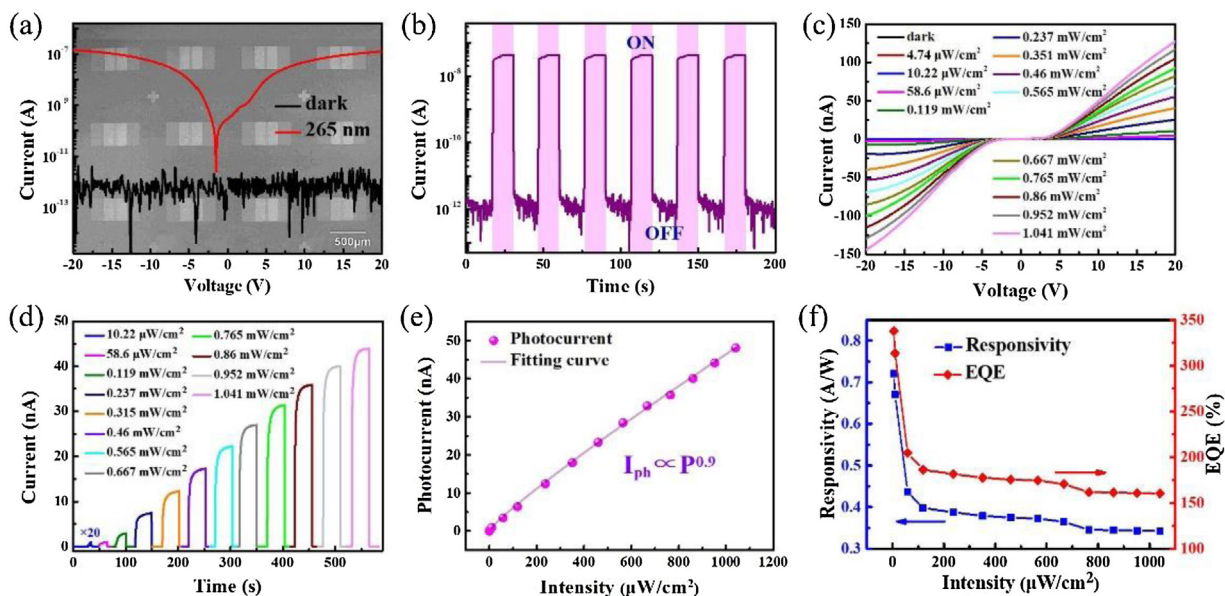
was found. Additionally, the optical bandgap was deduced to be  $\sim$ 5.15 eV from the curve of  $(\alpha h\nu)^2 - h\nu$  (where  $\alpha$ ,  $h$  and  $\nu$  denote absorbance, Planck's constant and frequency, respectively), which agreed well the reported values in literatures [2,7]. Therefore, by using this method, high-quality  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> thin films can be successfully prepared.

With the patterned  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> thin films as the building blocks, a photodetectors array composed of  $8 \times 8$  devices pixels was constructed. Fig. 3(d) depicts a FESEM image of a representative photodetector with a typical MSM device geometry. The elemental mapping profiles shown in Fig. 3(e) and (f) revealed clear and regular distribution boundaries of Au and Ga elements. Some unwanted Ga element detected beyond the region of the  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> thin film is probably caused by the residues during the lift-off process, which could be avoided by optimizing the experimental conditions in future work.

The optoelectrical properties of the as-synthesized  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> thin films were then characterized in order to study their potential application in deep-ultraviolet detector. Fig. 4(a) compares the current-voltage ( $I$ - $V$ ) curves of a typical photodetector in dark and



**Fig. 3.** (a) EDS spectrum and chemical composition of both gallium and oxygen atoms, (b) XRD pattern and (c) absorption spectrum of the  $\beta$ - $\text{Ga}_2\text{O}_3$  thin film, the inset in (c) shows the plot of  $(\alpha h\nu)^2 - h\nu$  for deducing the optical bandgap of the  $\beta$ - $\text{Ga}_2\text{O}_3$  thin film. (d) FESEM image of a typical  $\beta$ - $\text{Ga}_2\text{O}_3$  thin film-based photodetector with a MSM geometry. The elemental mapping showing the distribution of (e) Au and (f) Ga elements.



**Fig. 4.** (a)  $I$ - $V$  curves of the  $\beta$ - $\text{Ga}_2\text{O}_3$  thin film-based photodetector in dark and under 265 nm illumination ( $0.952 \text{ mW cm}^{-2}$ ). The background is a FESEM image of the photodetectors array. (b) Time-dependent photoresponse of the device under 265 nm illumination. (c)  $I$ - $V$  curves and (d) time-dependent photoresponse of the device under 265 nm illumination with different light intensities. (e) Photocurrent and (f) responsivity and EQE of the device as functions of light intensity.

under 265 nm DUV irradiation with a light intensity of  $0.952 \text{ mW cm}^{-2}$ . Evidently, the device had an extremely low dark current on the order of only pA even at a bias voltage of 20 V. It should be noted that the dark current was in fact over-valuated due to the limitations of our measurement conditions and instruments, considering that the measured dark current kept almost invariant within a large operational bias voltage range from -20 to 20 V in spite of some fluctuations. Interestingly, the channel current increased drastically by more than 4 orders of magnitude at 20 V when shinned by the DUV illumination, giving rise to a pronounced photoresponse. Furthermore, the time-dependent photoresponse was studied under the periodically turned DUV irradiation. As displayed in Fig. 4(b), the channel current could be readily switched between low- and high-

value states with excellent repeatability, indicating a good optical switching characteristic. At a bias of 10 V, the average values of dark current and photocurrent were  $\sim 6.2 \times 10^{-13} \text{ A}$  and  $\sim 3.8 \times 10^{-8} \text{ A}$ , respectively, which rendered a large  $I_{\text{light}}/I_{\text{dark}}$  ratio as high as  $\sim 6.13 \times 10^4$ .

The above photoresponse depended significantly on the incident DUV light intensity. Fig. 4(c) demonstrates the  $I$ - $V$  characteristics upon the DUV illumination with different light intensities in the range of  $4.74 \mu\text{W cm}^{-2}$  -  $0.565 \text{ mW cm}^{-2}$ . Apparently, the current rose gradually at both positive and negative bias voltages as the light intensity increased. Further study of time-dependent photoresponse under periodically switched DUV illumination with varied intensities also revealed an identical evolution tendency

**Table 1**  
Comparison of some key performance parameters of our device and  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> thin film-based DUV photodetectors in literatures.

Method	Dark current (A)	$I_{\text{light}}/I_{\text{dark}}$ ratio	Responsivity (A W <sup>-1</sup> )	$D^*$ (Jones)	$\tau_r/\tau_d$ (s)	Ref.
PLD	$1.88 \times 10^{-7}$	$\sim 10$	7.1 (10 V)	–	–	[38]
magnetron sputtering	$7.63 \times 10^{-9}$	$1.08 \times 10^3$	2.6 (10 V)	$1.6 \times 10^{12}$	0.26/1	[37]
magnetron sputtering	$1.0 \times 10^{-11}$	$\sim 10^5$	0.89 (10 V)	–	0.3/0.25	[24]
PECVD <sup>a</sup>	$4.0 \times 10^{-10}$	37	$2.0 \times 10^{-4}$ (0 V)	$6.9 \times 10^9$	–	[20]
PECVD	$4.0 \times 10^{-10}$	$\sim 10^5$	1.2 (10 V)	$1.9 \times 10^{12}$	–	[25]
OFZ <sup>b</sup>	$2.3 \times 10^{-10}$	$\sim 10^2$	4 (40 V)	–	0.3/0.2	[40]
Thermal-assisted conversion	$6.2 \times 10^{-13}$	$6.13 \times 10^4$	0.72 (10 V)	$4.18 \times 10^{11}$	1.1/0.03	This work

<sup>a</sup>) Plasma-enhanced chemical vapor deposition; <sup>b</sup>) Optical float-zone.

(Fig. 4(d)). This kind of relationship was reasonable because more photoinduced electron-hole pairs would be generated under the light irradiation with a higher intensity, which contributed to a larger photocurrent. In addition, we also noted that the  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> thin film-based detector could exhibit good optical switching feature for all illuminating conditions (Fig. 4(d)), suggesting that it could operate properly upon the DUV illuminations in a broad intensity regime. To gain more understanding of the above relationship, the dependence of photocurrent at 10 V on the light intensity was then fitted by using an extensively employed power law:  $I_{\text{ph}} \propto P^\theta$ , where  $I_{\text{ph}}$  and  $P$  represent the net photocurrent ( $I_{\text{ph}} = I_{\text{light}} - I_{\text{dark}}$ ) and the light intensity, and the index  $\theta$  is an empirical value reflecting the photocarrier recombination activity [34]. By carefully fitting the curve, a  $\theta = 0.9$  was attained in the light intensity of  $10.22 \mu\text{W cm}^{-2}$  -  $1.041 \text{ mW cm}^{-2}$  (Fig. 4(e)). The slight deviation from the ideal value ( $\theta = 1$ ) suggested the existence of photocurrent recombination loss in the device, which was likely due to the presence of some trap states between the Fermi level and the conduction band edge [35].

In order to compare the performance of varied DUV photodetectors more equitably, we then calculated several pivotal performance parameters including responsivity ( $R$ ) and external quantum efficiency (EQE). These two parameters are usually described by the following formula: [36]

$$R = \frac{I_{\text{ph}}}{P_\lambda S} = \text{EQE} \left( \frac{e\lambda}{hc} \right) \quad (1)$$

where  $P_\lambda$ ,  $S$ ,  $e$ ,  $\lambda$ ,  $h$ , and  $c$  denote the light intensity, the effective illumination area ( $S = 1.35 \times 10^{-4} \text{ cm}^2$ , the  $S$  equals to the channel area, which is determined by multiplying the channel length ( $\sim 45 \mu\text{m}$ ) by the channel width ( $\sim 300 \mu\text{m}$ ), the basic electronic charge ( $1.6 \times 10^{-19} \text{ C}$ ), the incident light wavelength, the Planck's constant and the speed of light, respectively. Therefore, the values of  $R$  and EQE were determined to be  $\sim 0.72 \text{ A W}^{-1}$  and  $\sim 338 \%$ , respectively, at the light intensity of  $4.74 \mu\text{W cm}^{-2}$ . Although the  $R$  value was slightly lower than that of DUV photodetectors based on  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> thin films prepared via physical vapor deposition (PVD) methods including magnetron sputtering and pulsed laser deposition (PLD) ( $0.89$ – $7.1 \text{ A W}^{-1}$ ) [24,37,38], it was still higher than that of devices comprising  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> thin films grown by MOCVD and low pressure chemical vapor deposition (LPCVD) techniques ( $0.32 \text{ mA W}^{-1}$ – $0.14 \text{ A W}^{-1}$ ) [9,39]. The  $R$  and EQE could be further improved by optimizing the device geometry, for example, shrinking the device channel width or employing interdigitated electrodes. In addition, it was found that both  $R$  and EQE values decreased monotonously with the increase in the light intensity. The results again verified that photocurrent recombination loss presented in the current device and could be ascribed to the intensified activity of photocarrier recombination due to increased density of charge carriers in the thin film when shinned by light illumination with a higher intensity.

Besides, specific detectivity ( $D^*$ ) is also one of the key figure-of-merits of a photodetector, which is usually employed to describe

the smallest detectable signal. This parameter can be expressed by the following equations [2,16]:

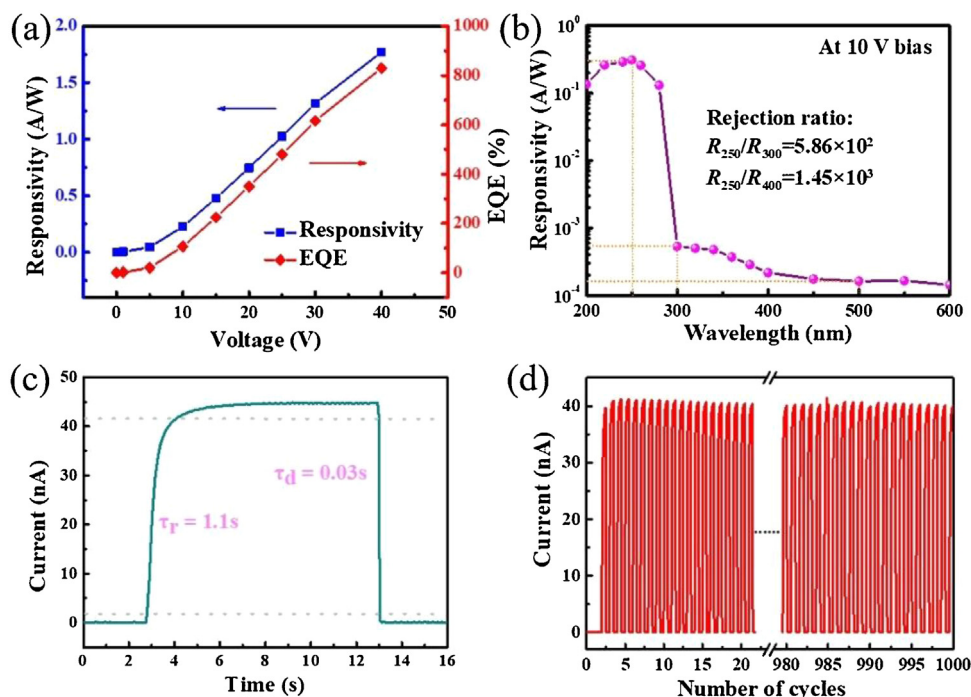
$$D^* = \frac{(AB)^{1/2}}{\text{NEP}} \quad (2)$$

$$\text{NEP} = \frac{i_n^{2/2}}{R} \quad (3)$$

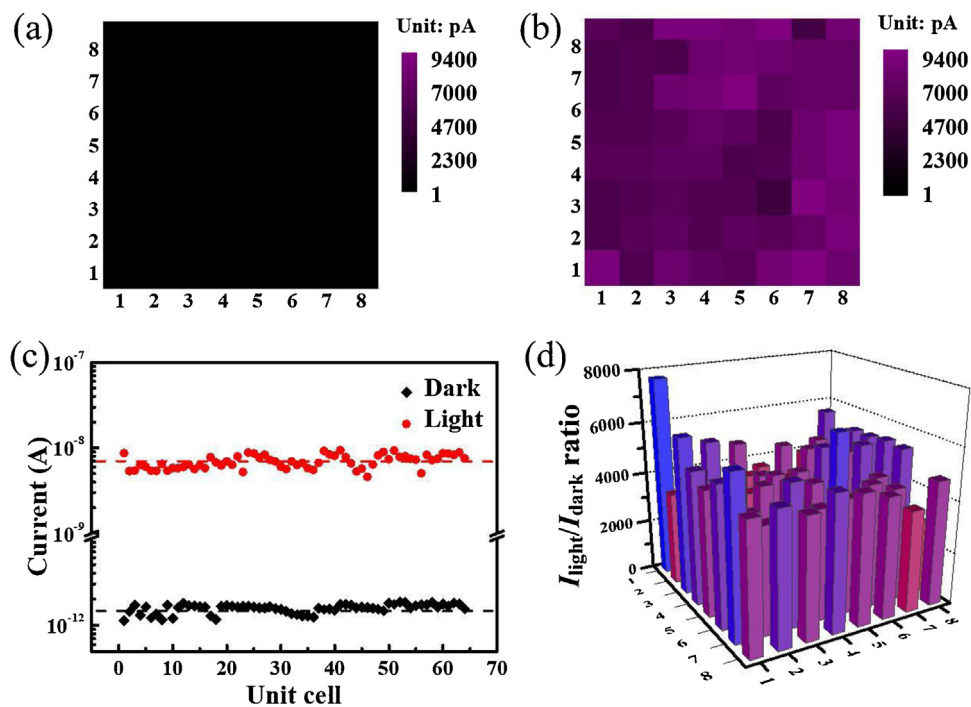
where  $A$ ,  $B$ ,  $\text{NEP}$  and  $i_n^{2/2}$  denote the effective area ( $A = 1.35 \times 10^{-4} \text{ cm}^2$ ), bandwidth, the noise equivalent power and the root-mean-square value of the noise current, respectively. The noise level at the bandwidth of 1 Hz was deduced to be  $\sim 2 \times 10^{-14} \text{ A Hz}^{-1/2}$  by performing Fourier transformation of the dark current recorded at 10 V bias (Fig. S3, in Supporting Information). Therefore, the  $\text{NEP}$  was  $\sim 2.78 \times 10^{-14} \text{ W Hz}^{-1/2}$ . Accordingly, the value of  $D^*$  at 10 V bias was calculated to be  $\sim 4.18 \times 10^{11}$  Jones. Note that such a decent  $D^*$  value was comparable with that of many previously reported  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> thin film-based DUV photodetectors [20,25,37]. Table 1 summarizes some important performance parameters of  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> thin film-based DUV photodetectors with MSM structure in literatures.

In addition to the incident light intensity, the operational bias voltage also had a great influence on the values of  $R$  and EQE. As shown in Fig. 5(a), both values rose progressively as the bias voltage increased. This was understandable because a larger external bias was helpful for accelerating separation of photocarriers and transfer of charge carriers, which would give a larger photoreponse. Further, the spectral response of our photodetector was studied, as plotted in Fig. 5(b). Evidently, the device exhibited a sizeable photoresponse in the DUV wavelength region from 200 to 280 nm with the peak response situated at  $\sim 250 \text{ nm}$ . As the light wavelength increased, the photoresponse declined dramatically, rendering a sharp response cut-off edge at  $\sim 280 \text{ nm}$ , which coincided with the optical absorption of the  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> thin films (Fig. 3(c)). What is more, the UV/visible rejection ratio defined as the ratio of responsivity at 250 and 400 nm ( $R_{250}/R_{400}$ ), and the solar-blind rejection ratio defined as the ratio of responsivity at 250 and 300 nm ( $R_{250}/R_{300}$ ) here were estimated to be  $\sim 1.45 \times 10^3$  and  $\sim 5.86 \times 10^2$ , respectively. The results implied that the present  $\beta$ -Ga<sub>2</sub>O<sub>3</sub> thin film-based DUV photodetector possessed an excellent solar-blind spectral selectivity.

The response speed that can reflect the capability of a photodetector to track quickly modulated photon signal was characterized as well. Fig. 5(c) displays a magnified photoresponse curve extracted from the time-dependent photoresponse, from which the rise and decay times ( $\tau_r$  and  $\tau_d$ ) were deduced to be  $\sim 1.1 \text{ s}$  and  $0.03 \text{ s}$ , respectively. Additionally, our photodetector also exhibited good reproducibility and stability. As depicted in Fig. 5(d), the time-dependent photoreponse could maintain almost unchanged with a good optical switching property even after thousands of cycles of operation. The slight fluctuation in photocurrent for each cycle was likely due to the disturbance existed in the measurement conditions. Furthermore, the air stability of our device was studied by re-measuring the time-dependent photoresponse after storage



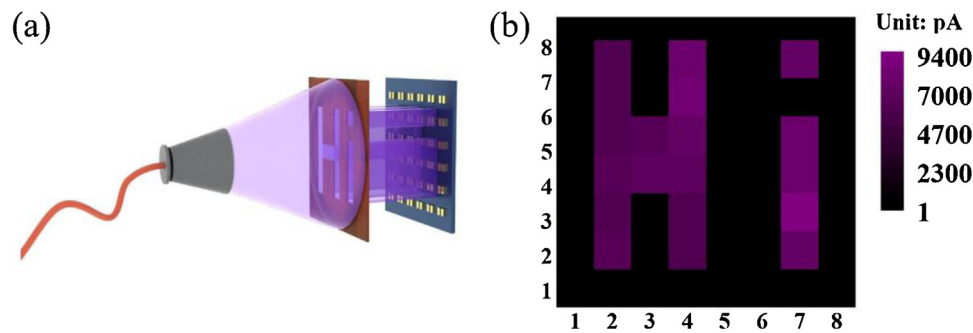
**Fig. 5.** (a) Responsivity and EQE of the device as functions of operational bias voltage. (b) Spectral responsivity of the device in the wavelength region of 200–600 nm. (c) An enlarged photoresponse curve shows the rise ( $\tau_r$ ) and decay times ( $\tau_d$ ). (d) Time-dependent photoresponse of the device over 1000 cycles of operation (265 nm,  $0.952 \text{ mW cm}^{-2}$ ).



**Fig. 6.** 2D current contrast maps showing the channel currents of all the  $8 \times 8$  devices (a) in dark and (b) upon 265 nm light illumination ( $0.25 \text{ } \mu\text{W cm}^{-2}$ ). (c) Channel current in dark and under 265 nm light irradiation for each device unit, the dash lines represent the average values for dark current (black) and photocurrent (red). (d) 3D diagram of the  $I_{\text{light}}/I_{\text{dark}}$  ratio for each device unit.

under ambient conditions without any protection for one month. As displayed in Fig. S4 (in Supporting Information), the device could keep its switching property and the photocurrent remained almost invariant, suggesting excellent air stability of the  $\beta\text{-Ga}_2\text{O}_3$  thin film-based DUV photodetector.

In order to explore the possibility of the current  $\beta\text{-Ga}_2\text{O}_3$  thin film-based photodetectors array for device integration application, the uniformity of the DUV photoresponse performance of all  $8 \times 8$  device units was further probed. Here, the dark current and photocurrent under a homogenous 265 nm DUV irradiation ( $0.25$



**Fig. 7.** (a) Schematic illustration of the setup for DUV light imaging. (b) The 2D current contrast map of the photodetectors array under 265 nm light illumination, showing the DUV light imaging function.

$\mu\text{W cm}^{-2}$ ) of each device were measured separately. For ease of comparison, the results were incorporated into 2D current contrast maps as demonstrated in Fig. 6(a) and (b), and the detailed data for all devices were provided in Fig. 6(c) as well. Obviously, these devices could operate properly. Careful analysis found that all device units had dark current in the range of 1–2 pA with an average value of  $\sim 1.57$  pA, and photocurrent in the range of 5–9 nA with an average values of  $\sim 5.98$  nA, at a bias of 10 V. The narrow fluctuations in both dark current and photocurrent suggested a uniform photoresponse performance. In addition, the  $I_{\text{light}}/I_{\text{dark}}$  ratio of all device units was extracted, which also showed a small fluctuation with the majority of the values in the range of  $3 \times 10^3$ – $6 \times 10^3$ . Such a good uniformity of the device performance could be attributed to the small variation in the thickness of the patterned  $\beta\text{-Ga}_2\text{O}_3$  thin films (Fig. S5 in Supporting Information), and also suggested a great potential for the present  $8 \times 8$  device units for integrated device applications.

Finally, the ability of the  $\beta\text{-Ga}_2\text{O}_3$  thin film-based photodetectors array to function as a DUV light image sensor was explored. As schematically shown in Fig. 7(a), a lab-built shadow mask with a pattern of “Hi” characters was placed between a homogenous 265 nm DUV illumination and the photodetectors array placed on a probe station. Only light illumination penetrating the hollow regime of the mask could shine on the devices while the rest of regimes were maintained in dark condition. While keeping the testing system still, the current of each device unit was recorded unit-by-unit by travelling a pair of probes connecting with a semiconductor parameter analyzer. After measurement of all units, the results were described by a 2D current contrast map. Significantly, as observed in Fig. 7(b), the shape of “Hi” could be clearly identified, which implied the preliminary DUV light imaging function of the  $\beta\text{-Ga}_2\text{O}_3$  thin film-based photodetectors array. The spatial resolution could be further improved by means of increasing the number of the device units as well as reducing the size of each device unit.

## Conclusion

In summary, we have demonstrated the periodic fabrication of high-quality  $\beta\text{-Ga}_2\text{O}_3$  thin films by using a facile patterned technique. Especially, benefiting from the technique flexibility of standard photolithography, such a method could be applied to prepare patterned  $\beta\text{-Ga}_2\text{O}_3$  thin films with arbitrary shapes in principle. With this technique, a  $\beta\text{-Ga}_2\text{O}_3$ -based photodetectors array comprising  $8 \times 8$  device units was constructed, which showed excellent DUV photoresponse performance, in terms of a large  $I_{\text{light}}/I_{\text{dark}}$  ratio exceeding  $10^4$ , a decent responsivity of  $\sim 0.72$  A  $\text{W}^{-1}$ , and a specific detectivity of  $\sim 4.18 \times 10^{11}$  Jones, respectively. In addition, the device could attain a large DUV/visible (250/400 nm) rejection ratio exceeding  $10^3$  with a steep response cut-off wavelength located at  $\sim 280$  nm. What is more, the excellent uni-

formity and repeatability of the performance of the device units allowed the present photodetectors array to function as a DUV light image sensor with a reasonable spatial resolution. Given the ease of processing, low-cost and good compatibility with the well-developed techniques, the proposed method was believed to have a huge potential for  $\beta\text{-Ga}_2\text{O}_3$ -based integrated optoelectronic devices with multifunctional applications.

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## Appendix A. Supplementary data

Supplementary material related to this article can be found, in the online version, at doi:<https://doi.org/10.1016/j.jmst.2020.09.015>.

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