

# Fast, Self-Driven, Air-Stable, and Broadband Photodetector Based on Vertically Aligned PtSe<sub>2</sub>/GaAs Heterojunction

Long-Hui Zeng, Sheng-Huang Lin, Zhong-Jun Li, Zhi-Xiang Zhang, Teng-Fei Zhang, Chao Xie, Chun-Hin Mak, Yang Chai, Shu Ping Lau, Lin-Bao Luo,\* and Yuen Hong Tsang\*

Group-10 layered transitional metal dichalcogenides including PtS<sub>2</sub>, PtSe<sub>2</sub>, and PtTe<sub>2</sub> are excellent potential candidates for optoelectronic devices due to their unique properties such as high carrier mobility, tunable bandgap, stability, and flexibility. Large-area platinum diselenide (PtSe<sub>2</sub>) with semiconducting characteristics is far scarcely investigated. Here, the development of a high-performance photodetector based on vertically aligned PtSe<sub>2</sub>-GaAs heterojunction which exhibits a broadband sensitivity from deep ultraviolet to near-infrared light, with peak sensitivity from 650 to 810 nm, is reported. The  $I_{\text{light}}/I_{\text{dark}}$  ratio and responsivity of photodetector are  $3 \times 10^4$  and 262 mA W<sup>-1</sup> measured at 808 nm under zero bias voltage. The response speed of  $\tau_r/\tau_f$  is 5.5/6.5  $\mu\text{s}$ , which represents the best result achieved for Group-10 TMDs based optoelectronic device thus far. According to first-principle density functional theory, the broad photoresponse ranging from visible to near-infrared region is associated with the semiconducting characteristics of PtSe<sub>2</sub> which has interstitial Se atoms within the PtSe<sub>2</sub> layers. It is also revealed that the PtSe<sub>2</sub>/GaAs photodetector does not exhibit performance degradation after six weeks in air. The generality of the above good results suggests that the vertically aligned PtSe<sub>2</sub> is an ideal material for high-performance optoelectronic systems in the future.

## 1. Introduction

The layered 2D transition disulfides (2D TMDs) with common formula of MX<sub>2</sub> (the M denotes a transition metal and X is a chalcogen atom, with strong in-plane covalent bonding and anisotropic layer structure with weak van der Waals bonding between layers), have attracted great attention in the past decade due to their distinctive material properties, for example, tunable open bandgap, high carrier mobility, and strong non-linear optical properties,<sup>[1–5]</sup> which may potentially bring the revolutionary changes in diverse fields including electronics,

photronics,<sup>[6,7]</sup> catalysis,<sup>[8]</sup> laser,<sup>[9]</sup> renewable energy,<sup>[10]</sup> and so on. As a new Group-10 TMDs based materials, PtSe<sub>2</sub> has come under spotlight for its tunable bandgap transition from monolayer (1.2 eV) to semimetal at bulk with zero bandgap. To date, a number of synthetic strategies have been developed for the fabrication of PtSe<sub>2</sub> layer. For instance, PtSe<sub>2</sub> flakes can be mechanically exfoliated from bulk PtSe<sub>2</sub> crystal but with very small and uncontrollable sample size.<sup>[11]</sup> To obtain larger size sample, Wang et al. for the first time grew monolayer PtSe<sub>2</sub> film through molecular beam epitaxy-assisted selenization process.<sup>[12]</sup> Recently, the efforts of synthesizing polycrystalline PtSe<sub>2</sub> film by either chemical reaction or thermally assisted conversion (TAC) at low temperature has been reported.<sup>[13–15]</sup> The PtSe<sub>2</sub> films derived by the above method are horizontally aligned. It has been experimentally observed that PtSe<sub>2</sub> exhibits very good stability in air (>1 year).<sup>[11,12]</sup> In addition, high mobility of PtSe<sub>2</sub> has been

predicted among the highest in TMDCs, which is comparable to that of black phosphorus.<sup>[11]</sup> Because of these appealing properties, PtSe<sub>2</sub> has shown potentially important applications in various nanodevices, such as field-effect transistors<sup>[16]</sup> and photocatalyst.<sup>[17]</sup> What's more, for the first time, a Si-PtSe<sub>2</sub> hybrid structure by combining vertically stacked PtSe<sub>2</sub> thin film with Si was fabricated, which can act as a sensitive gas sensor with extremely short response.<sup>[13]</sup> It was also found that like traditional metal–semiconductor (M–S) photodiode, the as-fabricated PtSe<sub>2</sub>-Si structure exhibited pronounced photovoltaic characteristics, and was highly sensitive to 920 nm laser illuminations with very high responsivity.

Near-infrared (NIR) light photodetector is vitally important for both military and civil applications including optical telecommunication, target imaging, remote version, time surveillance, and environmental monitoring.<sup>[18,19]</sup> Nowadays, the commercial IR photodetectors are dominated by narrow-gap semiconductors such as PbS, PbSe, PbTe, HgTe, and HgCdTe.<sup>[20,21]</sup> These materials normally entail the usage of sophisticated fabrication facility.<sup>[21–23]</sup> Silicon-based NIR photodetectors are also emerging recently. However, the fabrication of silicon NIR photodetectors is usually characterized by growth of narrow-gap Ge, GeSi, and HgCdTe thin film on Si using molecular

L.-H. Zeng, Dr. S.-H. Lin, C.-H. Mak, Prof. Y. Chai, Prof. S. P. Lau, Prof. Y. H. Tsang  
Department of Applied Physics and Materials Research Center  
The Hong Kong Polytechnic University  
Hung Hom, Kowloon, Hong Kong 99077, China  
E-mail: Yuen.Tsang@polyu.edu.hk

Prof. Z.-J. Li, Z.-X. Zhang, T.-F. Zhang, Prof. C. Xie, Prof. L.-B. Luo  
School of Electronic Science and Applied Physics  
Hefei University of Technology  
Hefei, Anhui 230009, China  
E-mail: luolb@hfut.edu.cn

DOI: 10.1002/adfm.201705970

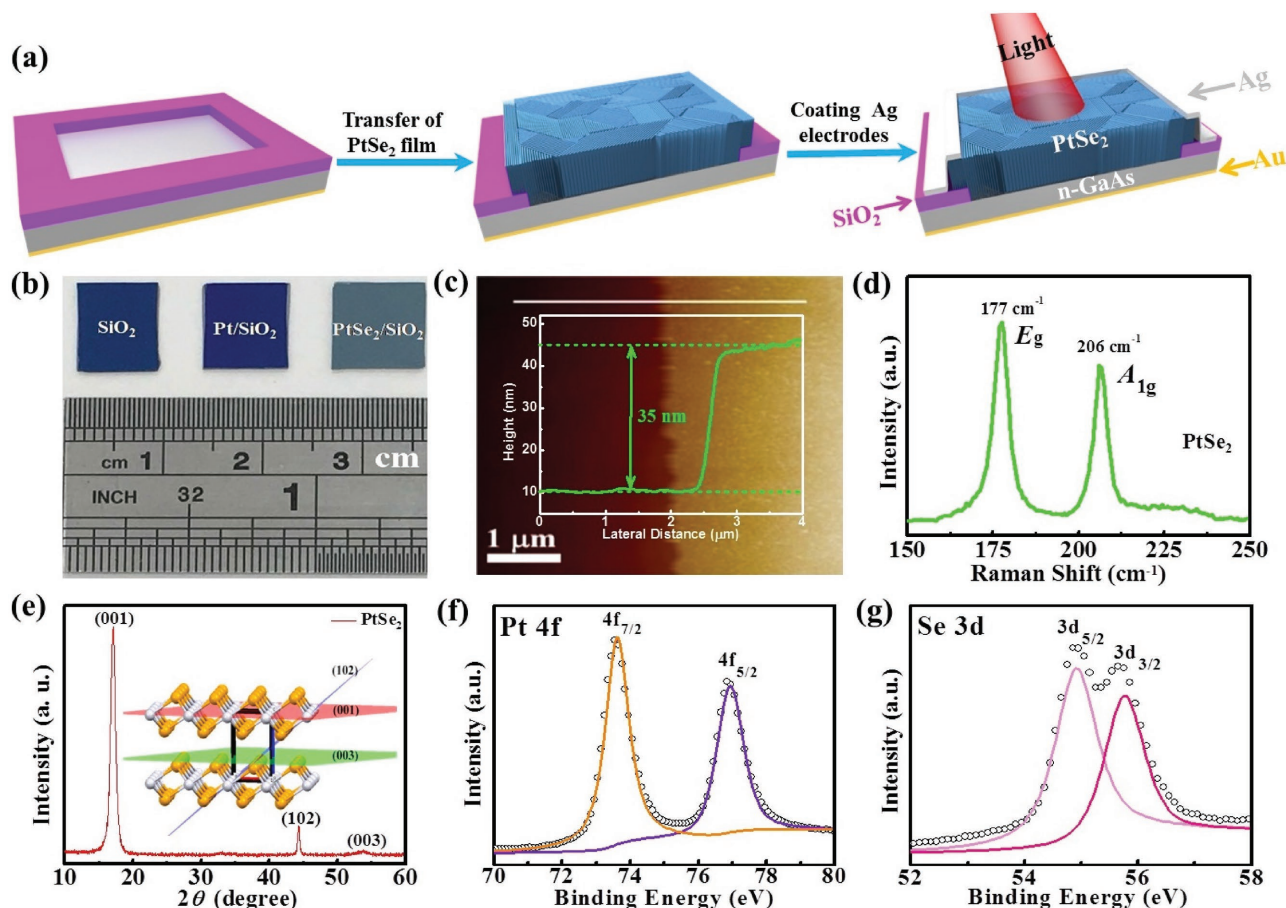
beam epitaxy (MBE). Even though the above devices have high sensitivity, it is undeniable that the fabrication process is time consuming. What's more, the fabrication cost is relatively high.<sup>[24,25]</sup> Considering the relatively small bandgap, PtSe<sub>2</sub> may find possible application in NIR detection.

In this work, we report on the synthesis of vertically aligned PtSe<sub>2</sub> multilayer film with semiconducting characteristics for high-performance photodetector application. The PtSe<sub>2</sub> films were synthesized by a simple selenization method, which has the obvious advantages in terms of low cost, large size, and well-controlled overall film thickness, in comparison with conventional chemical vapor deposition (CVD) method and mechanical exfoliation.<sup>[26]</sup> Based on the vertically aligned PtSe<sub>2</sub> multilayer film, we assembled a PtSe<sub>2</sub>-GaAs heterojunction which exhibited a broadband sensitivity to illumination ranging from deep ultraviolet (DUV) to near-infrared (NIR) light. Optoelectronic analysis revealed that the  $I_{on}/I_{off}$  ratio, responsivity, specific detectivity, and response speed of photodetector were  $3 \times 10^4$ , 262 mA W<sup>-1</sup>,  $2.52 \times 10^{12}$  Jones, and 5.5/6.5  $\mu$ s, respectively, which were comparable to or better than those of other TMDs photodetectors. Moreover, the present PtSe<sub>2</sub>/GaAs heterojunction showed obvious sensitivity in the range of visible-near-infrared, which was probably associated with the

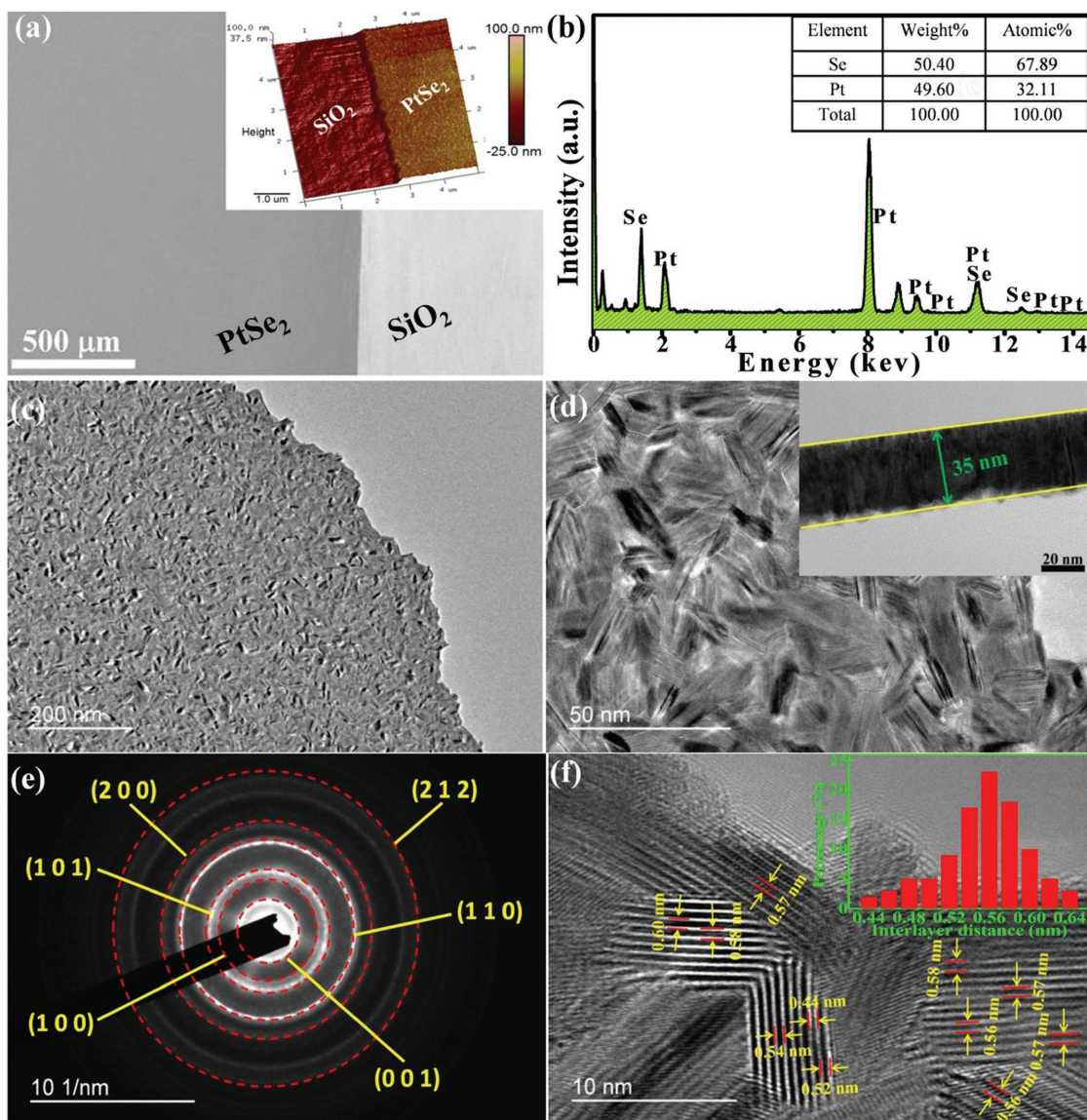
interstitial Se atoms, according to our DFT theoretical simulation. This study suggests that the present vertically aligned PtSe<sub>2</sub> film is a good candidate for assembling high-performance optoelectronic devices.

## 2. Results and Discussion

The scheme in Figure 1a illustrates the procedures to fabricate the PtSe<sub>2</sub>/GaAs heterojunction photodetector. The as-prepared PtSe<sub>2</sub> film was transferred onto the n-GaAs with a predefined window. An Ag paste was used as the top contact to PtSe<sub>2</sub>, whereas 50 nm Au served as the back contact to n-GaAs. In this study, the PtSe<sub>2</sub> was fabricated by a simple selenization of Pt thin film on a SiO<sub>2</sub>/Si wafer (Figure S1, Supporting Information). From the photograph shown in Figure 1b, it can be clearly seen that the color of Si/SiO<sub>2</sub> wafer changes slightly from navy to slate gray, when the Pt was transformed to PtSe<sub>2</sub>. Figure 1c displays an atomic force microscopy (AFM) image of the PtSe<sub>2</sub> film, in which a height profile along with the marked line is given in the inset diagram and the thickness of PtSe<sub>2</sub> is  $\approx$ 35 nm. With the availability of large-area PtSe<sub>2</sub> nanofilm, we are able to study the Raman spectrum of PtSe<sub>2</sub>. As revealed



**Figure 1.** a) Schematic illustration of the fabrication steps of PtSe<sub>2</sub>/GaAs heterojunction based photodetector. b) Photographs of SiO<sub>2</sub>/Si wafer, Pt/SiO<sub>2</sub>/Si, and PtSe<sub>2</sub>/SiO<sub>2</sub>/Si. c) AFM image of the PtSe<sub>2</sub> film on the SiO<sub>2</sub>/Si substrate, the inset shows the height profile along the white line marked in (c). d) Raman spectra of PtSe<sub>2</sub> nanofilm. e) XRD pattern of PtSe<sub>2</sub> film and inset shows 3D atomic structure diagram of sample. f,g) The high-resolution XPS spectra of Pt 4f and Se 3d.



**Figure 2.** a) The FESEM image of PtSe<sub>2</sub> image on a SiO<sub>2</sub>/Si wafer, the inset shows 3D AFM image of PtSe<sub>2</sub>. b) The EDS spectrum of the PtSe<sub>2</sub> sample. c,d) TEM images of synthesized PtSe<sub>2</sub> nanofilm with different magnification. e) The SAED pattern of the PtSe<sub>2</sub> sample. f) HRTEM image of the PtSe<sub>2</sub>, the inset shows statistical distribution of interlayer distance.

in Figure 1d, the peak centered at 177 cm<sup>-1</sup> can be assigned to the E<sub>g</sub> in-plane vibrational mode of Se atoms and the peak centered at 206 cm<sup>-1</sup> to the A<sub>1g</sub> out-of-plane vibration mode.<sup>[27]</sup> Further X-ray diffraction (XRD) pattern in Figure 1e is characterized by three diffraction peaks located at 17.14°, 44.28°, and 53.86° due to the (001), (102), and (003) crystal planes of PtSe<sub>2</sub>. According to the X-ray photoelectron spectroscopy (XPS) analysis, there are two peaks at 73.55 and 76.85 eV that are related to Pt 4f<sub>7/2</sub> and Pt 4f<sub>5/2</sub> orbitals, respectively, whereas the binding energies for the Se 3d<sub>5/2</sub> and Se 3d<sub>3/2</sub> are located at 54.85 and 55.65 eV.<sup>[16]</sup>

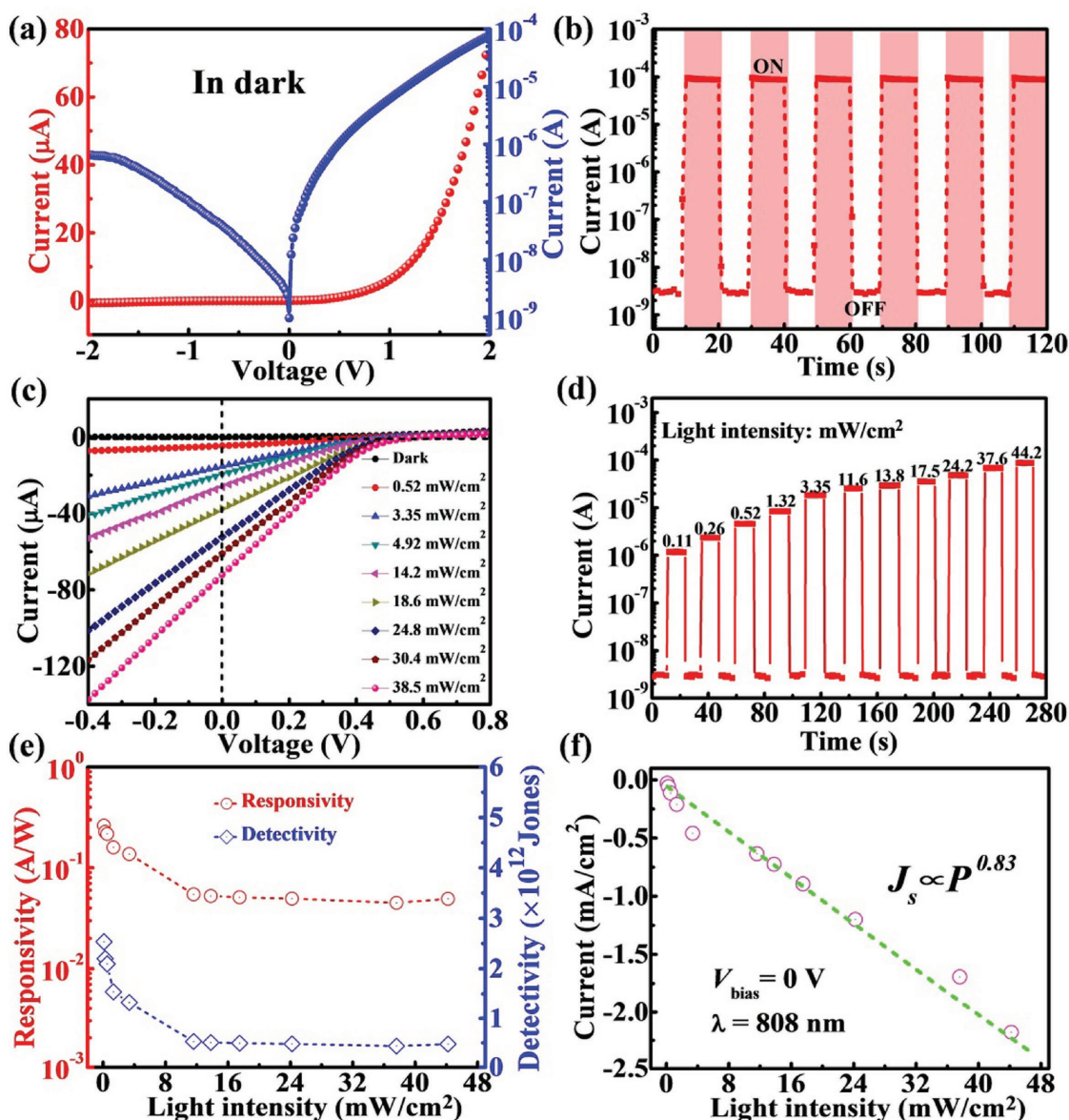
By studying the morphology using field emission scanning electron microscopy (FESEM), it can be easily found that the as-synthesized film is continuous and uniform in Figure 2a. The magnified 3D AFM surface topography was shown in the right

panel of Figure 2a, indicating the PtSe<sub>2</sub> film has a relatively smooth surface. The atomic ratio of Se/Pt inspected by energy-dispersive X-ray spectroscopy (EDS) was determined to be 2.11/1 (inset of Figure 2b), which is slightly higher than the stoichiometric composition of PtSe<sub>2</sub>. Figure 2c,d shows the transmission electron microscopy (TEM) images of the PtSe<sub>2</sub> layer at different magnifications, revealing that the PtSe<sub>2</sub> is polycrystalline with a vertically aligned layered structure. This special structure of PtSe<sub>2</sub> was different from the result of ref. [13], in which vertically stacked PtSe<sub>2</sub> was obtained.<sup>[13]</sup> Understandably, such a difference is related to the usage of Pt thin film precursors with different thickness (0.5 nm vs 12 nm). According to previous study, such a difference in Pt thickness means different density, compactness, and roughness of Pt, which therefore leads to PtSe<sub>2</sub> with different alignment during selenization

process.<sup>[28,29]</sup> The detailed reason for the formation of vertically aligned PtSe<sub>2</sub> is complicated and needs further investigation. The thickness of the PtSe<sub>2</sub> film is estimated to be 35 nm, according to cross-sectional TEM image as shown in the inset of Figure 2d. Further diffraction rings from inside to outside can be readily ascribed to the (001), (100), (101), (110), (200), and (212) planes of PtSe<sub>2</sub>.<sup>[30]</sup> It should be noted that the interlayer distance of the vertically grown PtSe<sub>2</sub> has a wide distribution from 0.44 to 0.64 nm, with an average value of 0.58 nm. Such an interlayer distance is relatively larger than previously reported value ( $\approx 0.52$  nm).<sup>[17]</sup> Considering the excessive amount of Se atoms in the PtSe<sub>2</sub> (Pt/Se = 1/2.11) measured by EDS, it can be easily concluded that such an increase in interlayer distance

is due to the interstitial Se atoms, which are normally located at gap between the interlayers. In addition, according to scanning transmission electron microscope (STEM) image and EDS elemental mapping analysis, both Pt and Se are homogeneously distributed over the whole film in Figure S2 (Supporting Information).

The as-fabricated PtSe<sub>2</sub>/GaAs heterojunction exhibits a typical rectifying behavior without light illumination, as indicated by the dark curve in Figure 3a. In light of the good contact of both Ag/PtSe<sub>2</sub> and Au/GaAs (Figure S3, Supporting Information), the above nonlinear  $I$ - $V$  characteristics stem from heterojunction formed at PtSe<sub>2</sub>/GaAs contact. The rectification ratio is determined to be  $\approx 120$  at  $\pm 2$  V, which is higher



**Figure 3.** a) The  $I$ - $V$  curves of the heterojunction measured under dark. b) Photoresponse of the device under 808 nm ( $44.2 \text{ mW cm}^{-2}$ ) light illumination at zero bias. c)  $I$ - $V$  curves of the device at varied light intensity illustration (808 nm). d) Photoresponse of the photodetector under various light intensity at 0 V. e) Both responsivity and specific detectivity of the PtSe<sub>2</sub>/GaAs heterojunction under 808 nm light illumination with different intensities. f) Photocurrent density as a function of the incident-light intensity at zero bias.

than previously studied graphene/GaAs Schottky junction ( $\approx 24$ ,  $\approx 49$ ,  $\approx 100$ ).<sup>[18,31,32]</sup> Figure 3b shows the photoresponse characteristics when the incident light (808 nm,  $44.2 \text{ mW cm}^{-2}$ ) was switched on and off repeatedly. Clearly, the current of heterojunction significantly increases with light irradiation on, giving stable and repeatable  $I_{\text{on}}/I_{\text{off}}$  of  $3 \times 10^4$ . It is also revealed that the fabricated device can be reversibly switched quickly between low and high current states with steep rise and fall edges, indicating that electron–hole pairs could be effectively generated and separated in the PtSe<sub>2</sub>/GaAs heterojunction. In fact, the photocurrent of the device is highly dependent on the light intensity. Figure 3c plots the  $I$ – $V$  curves under the light intensities ranging from  $0.52$  to  $38.5 \text{ mW cm}^{-2}$ . Apparently, with the increase of light intensity, the photocurrent is observed to increase gradually. So, it is easily noted that the device exhibits pronounced photovoltaic behaviors under light illumination within the measurement range from  $-0.4$  to  $+0.4 \text{ V}$ , indicating the PtSe<sub>2</sub>/GaAs heterojunction can function as a self-driven photodetector.<sup>[33]</sup> Further photoresponse characteristic in Figure 3d shows that the PtSe<sub>2</sub>/GaAs device can be switched between on and off states under varied light intensity with good reproducibility.

To evaluate the device performance of the current PtSe<sub>2</sub>/GaAs photodetector, the responsivity ( $R$ ), specific detectivity ( $D^*$ ), and response speed were then calculated.  $R$  is defined as the photocurrent generated per unit power of the incident light on the effective area and was obtained from the experimental data by using the formula

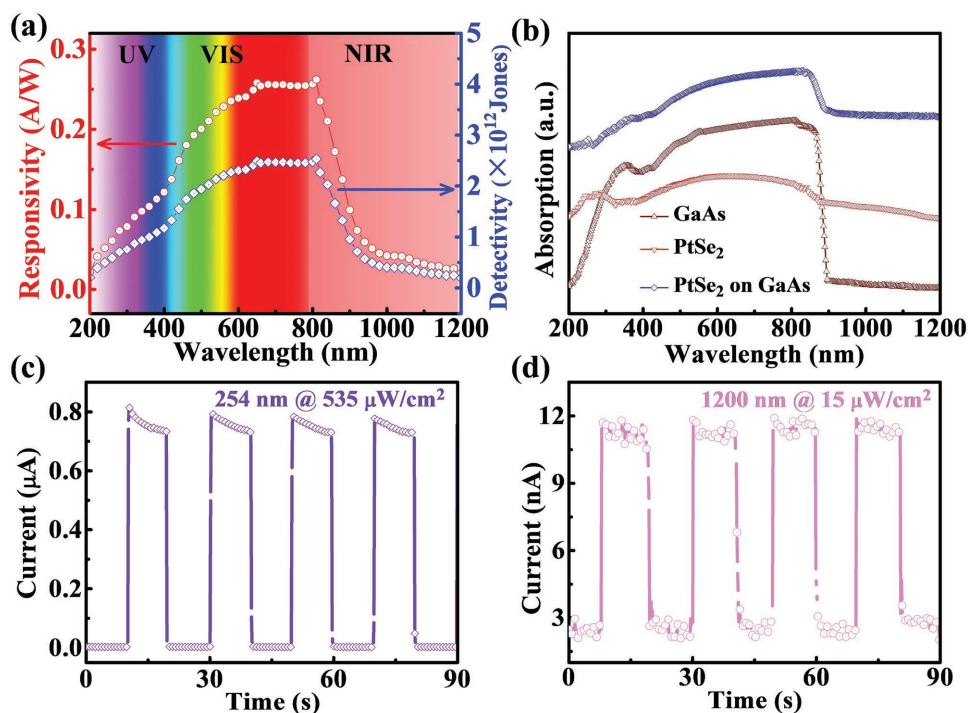
$$R(\text{A W}^{-1}) = \frac{I_p - I_d}{P_{\text{opt}} S} \quad (1)$$

where  $I_p$  is the photocurrent,  $I_d$  the dark current,  $P_{\text{opt}}$  the incident-light intensity, and  $S$  the effective illuminated area ( $S = 0.04 \text{ cm}^2$ ). On the other hand,  $D^*$  represents the ability of a detector to detect weak optical signals, which can be calculated from the following equation

$$D^* = A^{1/2} R / (2qI_d)^{1/2} \quad (2)$$

where  $A$  is the effective area of device ( $A = 0.09 \text{ cm}^2$ ),  $q$  the electronic charge, and  $R$  the responsivity, respectively. Based on Equations (1) and (2),  $R$  and  $D^*$  of PtSe<sub>2</sub>/GaAs heterojunction were estimated to be  $262 \text{ mA W}^{-1}$  and  $2.52 \times 10^{12} \text{ Jones}$  under light intensity of  $0.11 \text{ mW cm}^{-2}$  at zero bias voltage, respectively. Figure 3e plots both  $R$  and  $D^*$  under illumination of various light intensities. It is noted that both  $R$  and  $D^*$  of photodetector reduced with increasing light intensity. This finding is understandable as the photogenerated electrons can be captured by the trap states under lower light intensity, therefore leading to a reduction in recombination possibility.<sup>[34,35]</sup> In addition, the photocurrent current as a function of the light intensity is shown in Figure 3f and the curve was fitted to obtain law of  $J_s \propto P^{0.83}$ . The nonunity exponent of the law is due to the complex process of electron–hole generation, trapping, and recombination activity within the PtSe<sub>2</sub>/GaAs heterojunction.<sup>[36]</sup>

Next, the wavelength-dependent responsivity and specific detectivity of the PtSe<sub>2</sub>/GaAs device were studied. Figure 4a shows the continuous spectral responsivity in the range from 200 to 1200 nm. Obviously, this device exhibits broad photoresponse, with peak response in the range from 650 to 810 nm.



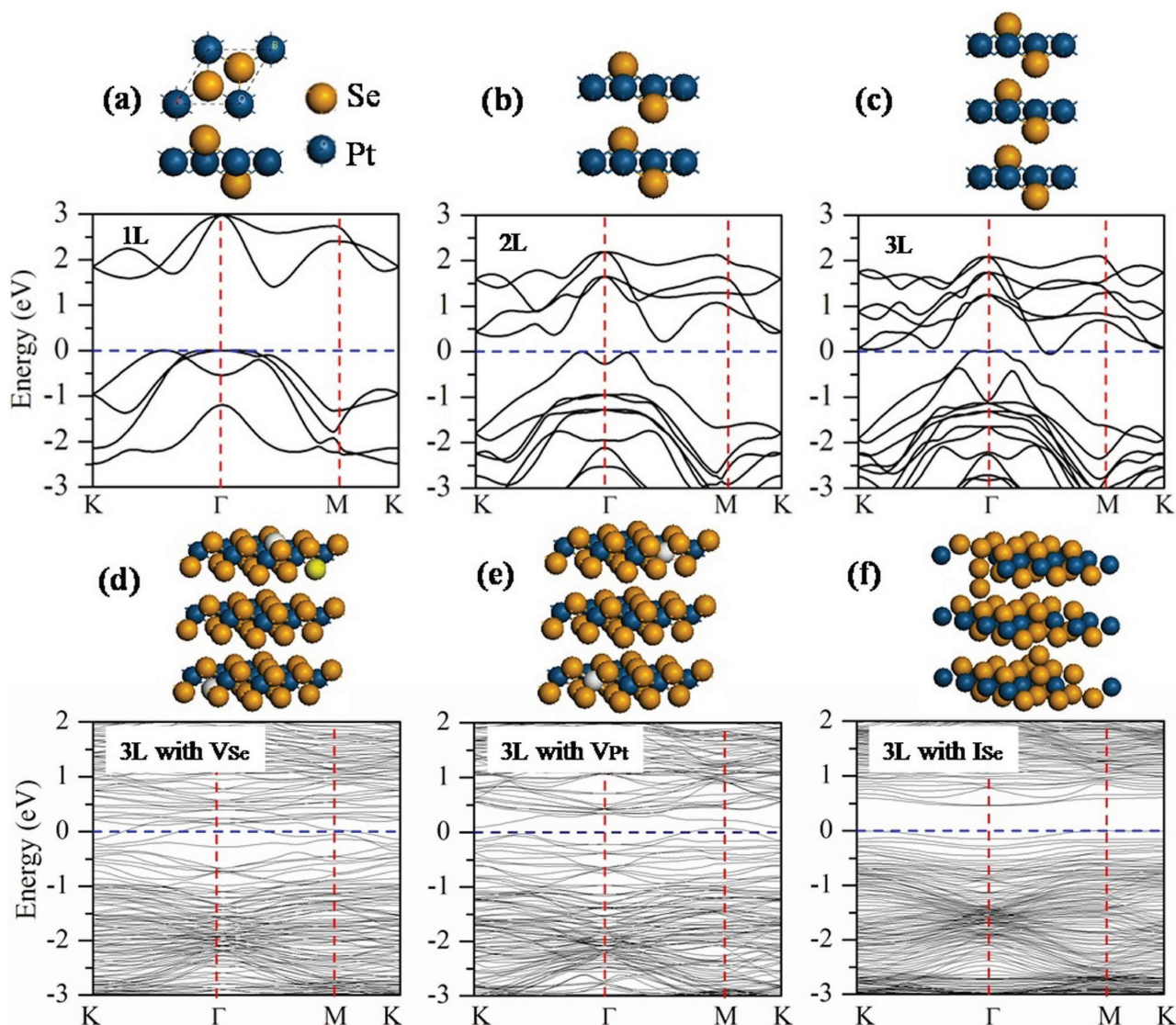
**Figure 4.** a) Wavelength-dependent responsivity and specific detectivity of PtSe<sub>2</sub>/GaAs. b) Absorption spectrum of the PtSe<sub>2</sub> film on GaAs substrate. The absorption spectra of PtSe<sub>2</sub> grown on quartz substrate under same conditions and bare substrate were also plotted for comparison. Time-dependent photoresponse of PtSe<sub>2</sub>/GaAs heterojunction excited by pulsed light at c) 254 and d) 1200 nm (d) under zero bias condition.

Such peak response can be exclusively attributed to the contribution from PtSe<sub>2</sub> layer, when compared with the spectral responsivity of PtSe<sub>2</sub>, GaAs, and Pt/GaAs devices, shown in Figure S4 (Supporting Information). In Figure S4a (Supporting Information), the pure PtSe<sub>2</sub> exhibits a high responsivity from 500 to 810 nm and the responsivity values decrease when the incident light wavelength outside this wavelength region. Due to the limited light absorption arising from the small thickness, the absolute responsivity value of the pure PtSe<sub>2</sub> is much lower than that of PtSe<sub>2</sub>/GaAs. On the other hand, for PtSe<sub>2</sub>/GaAs device, the photocurrent is mainly contributed by the generation of photocarriers in GaAs, which could be confirmed by the much higher responsivity value and the sharply decreased responsivity at ≈810 nm similar to that of pure GaAs and Pt/GaAs devices (Figure S4b,c, Supporting Information). In addition, due to the relatively strong light absorption of PtSe<sub>2</sub> from 500 to 810 nm, photocarriers generated in PtSe<sub>2</sub> also contribute to the photocurrent, which eventually gives a high responsivity from 500 to 810 nm with peak values from 650 to 810 nm. The maximum *R* and *D*\* values are ≈0.262 A W<sup>-1</sup> and 2.52 × 10<sup>12</sup> Jones, which are comparable to the reported values for graphene-based photodetector (*R* ≈ 1.73 mA W<sup>-1</sup>, *D*\* ≈ 1.83 × 10<sup>11</sup> Jones),<sup>[18]</sup> MoS<sub>2</sub>/GaAs (*R* ≈ 0.43 mA W<sup>-1</sup>, *D*\* ≈ 2.28 × 10<sup>11</sup> Jones),<sup>[37]</sup> and multilayer MoS<sub>2</sub> phototransistor (*R* ≈ 120 mA W<sup>-1</sup>, *D*\* ≈ 10<sup>10</sup> Jones).<sup>[38]</sup> As we will discuss later, this abnormal spectral selectivity is associated with the semiconducting characteristics of PtSe<sub>2</sub> layer. It is worth noting that even though the responsivity is relatively low in deep UV and near-infrared range, the present PtSe<sub>2</sub> films based device is still able to detect both deep UV and near-infrared light illumination. As shown in Figure 4c,d, the PtSe<sub>2</sub>/GaAs device can be readily switched between low-/high-conduction states, when the illumination repeatedly turn on and off. The excellent reproducibility suggests that this device may also find application for broadband photodetectors.<sup>[39]</sup>

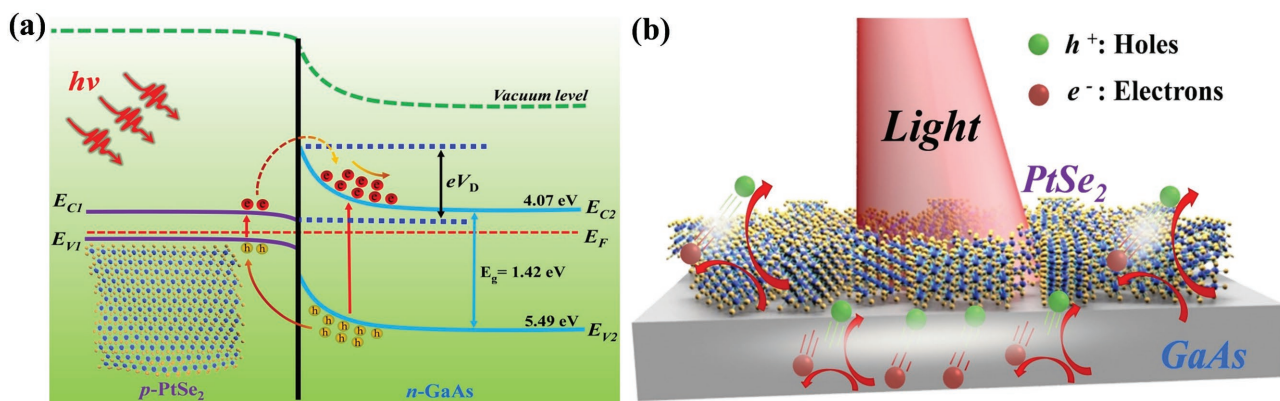
It is known that GaAs based photodetectors usually exhibit the strongest sensitivity at around 850 nm,<sup>[18,40]</sup> due to the band gap of 1.42 eV. This is quite different from the present GaAs/PtSe<sub>2</sub> device with peak photoresponse in the range from 650 to 810 nm. Such a difference in spectral selectivity is without question attributed to the PtSe<sub>2</sub>. To unveil the authentic origin behind this phenomenon, we then studied the electronic band of various structures by first-principle density functional theory (DFT) calculation. The geometrical and band structures of the PtSe<sub>2</sub> thin films with one, two, and three layers are presented in Figure 5a–c. One can see from the figure that the PtSe<sub>2</sub> thin films with one and two layers show semiconducting characteristics, and the bandgap clearly becomes smaller from one layer to two layers. As the thickness of thin films increase to three layers, the metallic character occurs. Therefore, following this decreasing trend of bandgap, one can conclude that the present freestanding PtSe<sub>2</sub> thin films possess metallic behavior, which is consistent with previous study.<sup>[11,41]</sup> Like graphene, this metallic behavior however cannot induce a change in spectral selectivity in that the bandgap was not opened up. So, the doping effects on electronic structures of the PtSe<sub>2</sub> thin films, Se-vacancy, Pt-vacancy, and the interstitial incorporation of Se between two neighboring layers are then taken into account in the supercells. In view of the

computational cost, the 3 × 3 supercell of the PtSe<sub>2</sub> thin films with three layers is employed in the calculation. In this supercell, the separation between the neighboring dopants is larger than 11 Å. This distance facilitates to reduce the interaction among the dopants. On the other hand, the PtSe<sub>2</sub> thin films with three layers have similar metallic character as the thin film with more layers. The Se-vacancy and Pt-vacancy are modeled in the topmost and bottommost layers, and the interstitial atoms of Se are incorporated between the two neighboring layers of the supercell. The band structures in Figure 5d–f indicate that the PtSe<sub>2</sub> thin films with the Se-vacancy and Pt-vacancy both possess metallic characters, whereas the system with the interstitial incorporation of Se is semiconducting. It is noted here that in the present work the PtSe<sub>2</sub> thin films are synthesized in the Se-rich condition. This will facilitate Se atoms to interstitially incorporate into parallel layers. In this type of doping, the atomic ratio between Se and Pt is more than 2, which are consistent with the present experimental measurements (see inset of Figure 2b). To further study the concentration influence of the interstitial incorporation of Se on electronic structures, the band structures of the 2 × 2 and 4 × 4 supercells with three layers are calculated and presented in Figure S5 (Supporting Information). It can be seen from the figure that the defect levels are largely determined by the doping concentration. Based on the above results, it can be easily concluded that the interstitial incorporation of Se between two layers (see Figure 2f), can lead to the opening of bandgap of PtSe<sub>2</sub>, and therefore will be responsible for the shift in spectral response.<sup>[42]</sup>

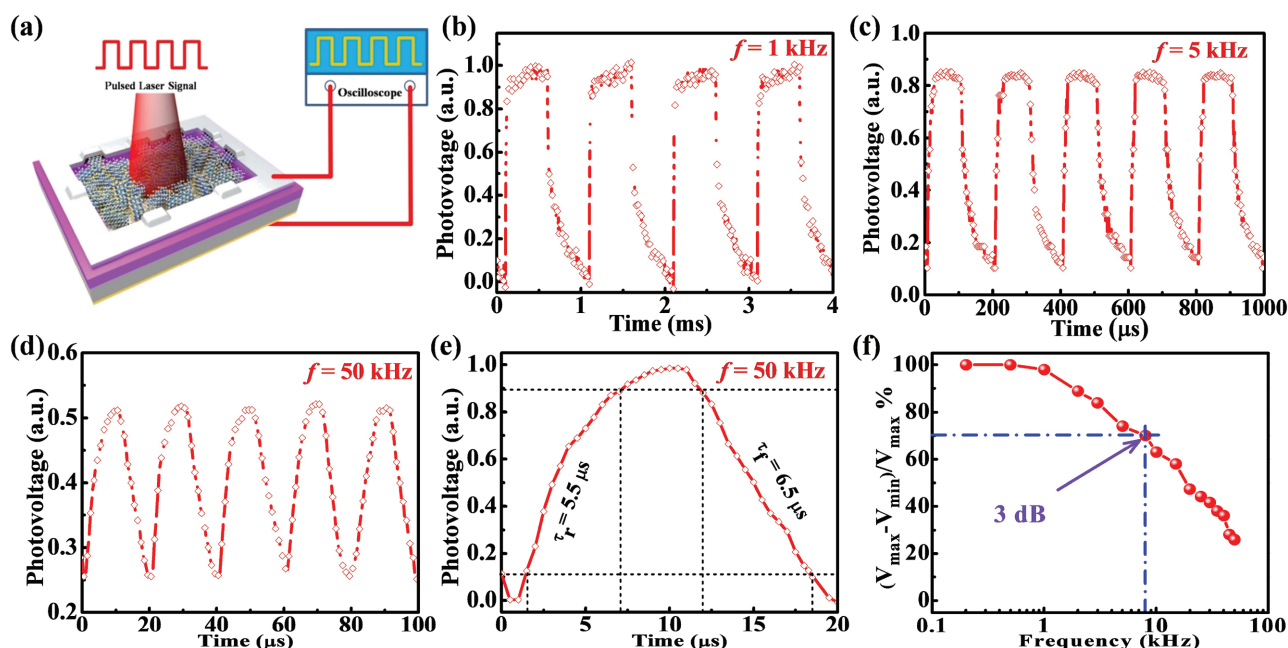
To further understand the photoresponse and spectral selectivity of the photodetector under light irradiation, we correlated the response characteristic with carrier transport process and the energy band diagrams of PtSe<sub>2</sub>/GaAs heterojunction structure and proposed a mechanism here based on the above results (see Figure 6a). Thus, work function of PtSe<sub>2</sub> is investigated by using ultraviolet photoemission spectroscopy (UPS) as shown in Figure S6a (Supporting Information), indicating that Fermi level of the PtSe<sub>2</sub> is located at 5.12 eV,<sup>[43]</sup> which is higher than that of GaAs. The typical transfer of the fabricated transistor based on PtSe<sub>2</sub> films is shown in Figure S6b (Supporting Information), indicating the p-type behavior of prepared PtSe<sub>2</sub> films. Due to difference in Fermi level (*E*<sub>F</sub>), once the p-type PtSe<sub>2</sub> is in contact with n-type GaAs, the electrons would diffuse from GaAs into PtSe<sub>2</sub> while holes should diffuse from PtSe<sub>2</sub> in to GaAs. Therefore, the energy levels near the GaAs surface will bend upward whereas the energy levels near PtSe<sub>2</sub> surface will bend downward, eventually the Fermi levels of PtSe<sub>2</sub> and GaAs align in the same level, leading to the formation of built-in electric field near the PtSe<sub>2</sub>/GaAs interface.<sup>[44]</sup> As often observed in graphene/semiconductor interface, under light irradiation, the electron–hole pairs will be generated in GaAs at the interface, and then separated by the built-in electric field as shown in Figure 6b, giving rise to generation of photocurrent. Due to the interstitial Se atoms during this photosensing process, parts of PtSe<sub>2</sub> layer will also strongly absorb incident light with wavelength of 650–810 nm, converting the photons to electron–hole pairs, and finally contribute to the photocurrent in the external circuit. Namely, the PtSe<sub>2</sub>/GaAs heterojunctions display the stronger sensitivity between 650 and 810 nm.



**Figure 5.** Geometrical and electronic band structures of the PtSe<sub>2</sub> primitive cells with a) one-layer, b) two-layer, c) three-layer, and the 3 × 3 supercells with d) Se-vacancy, e) Pt-vacancy, and f) interstitial incorporation of Se. The vacant atoms of Se and Pt are denoted by the white balls in (d) and (e). The Fermi level is located at 0 eV.



**Figure 6.** a) The energy band diagram of the PtSe<sub>2</sub>/GaAs heterojunction at zero bias under illumination. b) Schematic illustration describing the rapid separation and movement of photoexcited carriers under illumination.



**Figure 7.** a) The schematic illustration of the setup for studying the time response of the photodetector. Photoresponse of the self-driven PtSe<sub>2</sub>/GaAs heterojunction photodetector to pulsed light irradiation (808 nm) with a frequency of b) 1, c) 5, and d) 50 kHz, respectively. e) A single normalized cycle measured at 50 kHz for estimating both response time ( $\tau_r$ ) and recovery time ( $\tau_f$ ). f) Relative balance  $(V_{\max} - V_{\min})/V_{\max}$  versus switching frequency, showing the 3 dB cutoff frequency of  $\approx 8$  kHz.

The response speed is another critical parameter of photodetectors, which reflects the ability of a photodetector to follow a varied optical signal.<sup>[39,45]</sup> The response speed of PtSe<sub>2</sub>/GaAs heterojunction was measured by using a chopped laser beam with varied frequencies to generate pulsed light. The schematic illustration of the measurement is shown in **Figure 7a**, in which the temporal photoresponse signal (photovoltage with time) was recorded by a digital oscilloscope under the illumination of the pulsed NIR diode laser with several different frequencies from 1 kHz to 50 kHz. **Figure 7b–d** shows that the photoresponse of the heterojunction device to pulsed light. Obviously, the response is very fast with excellent good repeatability in the frequency 1, 5, and 50 kHz. From the magnified photoresponse curve in **Figure 7e**, the response time can be quantified. In the time domain, the speed of a photodetector is often characterized by the rise/fall time ( $\tau_r/\tau_f$ ). The time needed for the current to increase from 10% to 90% of the peak value or vice versa is defined as the rise and fall time, respectively.<sup>[46]</sup> Our device yield a rise time ( $\tau_r$ ) and fall time ( $\tau_f$ ) of 5.5 and 6.5  $\mu$ s at zero external bias voltage, respectively. However, this response speed is slightly slower than the results at  $-2$  V at the same frequency of 50 kHz (for optoelectronic study and test at  $-2$  V refer to **Figure S7**, Supporting Information) due to a reduced depletion layer capacitance under reverse bias.<sup>[47]</sup> It is noted that the response speed is faster than other 2D materials based photodetectors (see **Table 1**).<sup>[48–53]</sup> Such a relatively fast response speed is related to the special structure, which is beneficial for the separation and transport of photogenerated electron–hole carriers. As a matter of fact, this response speed can be further optimized by reducing the capacitance of the device or the thickness of the GaAs material. Besides, the relative balance  $(V_{\max} - V_{\min})/V_{\max}$  of the photovoltage as a function of fre-

quency in a range of 200 Hz to 50 kHz is depicted in **Figure 7f**. The 3 dB frequency that was defined as the frequency at which response dropped 0.707 of the initial value was estimated to be 8 kHz and the relative balance only decreases by less than 20% even at a high frequency of 50 kHz, suggesting that the PtSe<sub>2</sub>/GaAs heterojunction photodetector is capable of monitoring fast optical signals.<sup>[54]</sup>

The long-term stability of photodetector remains an important concern for device application. Therefore, we investigate the air stability of PtSe<sub>2</sub>/GaAs photodetector by placing sample in air for six weeks without any encapsulation. Remarkably, the device showed high stability, and no obvious degradation in sensitivity was observed (**Figure 8a,b**). The high device stability is not only directly associated with excellent air stability of PtSe<sub>2</sub> film and GaAs, but also the high quality of heterojunction based PtSe<sub>2</sub>/GaAs. As illustrated in **Figure 8c**, the XPS spectra of the aged PtSe<sub>2</sub> sample are virtually unchanged after six weeks storage in air. Moreover, the Raman spectra of the pristine PtSe<sub>2</sub> and the sample after six weeks aging in air are shown in **Figure 8d**. It is clearly observed that Raman spectra of the aged sample are similar to the pristine PtSe<sub>2</sub>. These results indicate the stability of PtSe<sub>2</sub> in air is much better than that of black phosphorus.<sup>[55]</sup>

### 3. Conclusion

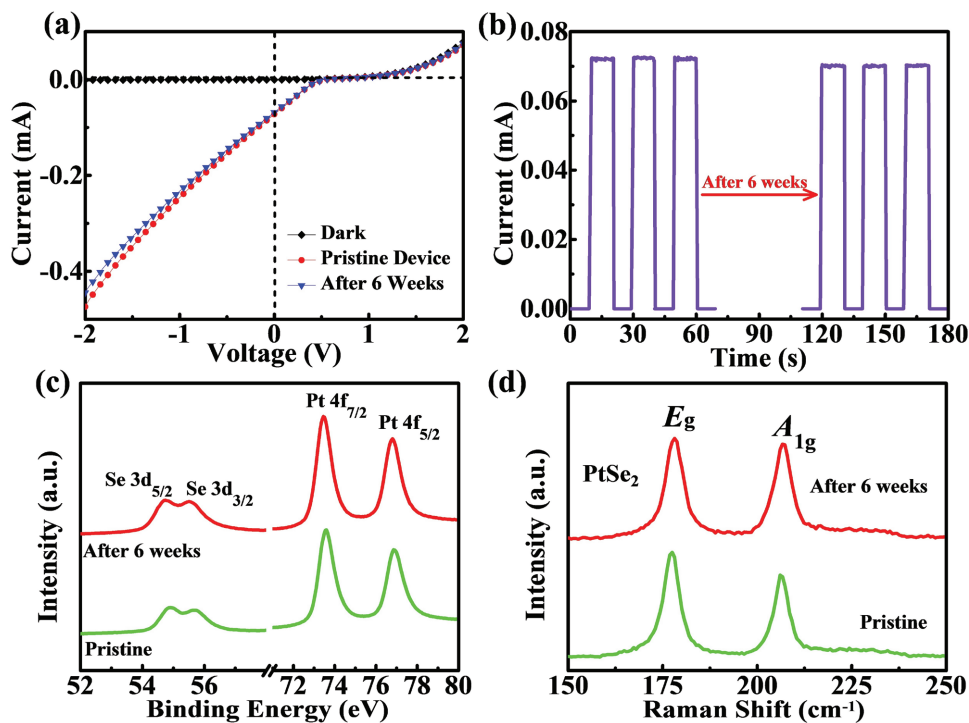
In summary, we have demonstrated an air-stable, self-powered, and broadband photodetector based on vertically aligned structure PtSe<sub>2</sub>/GaAs heterojunction. Such a heterojunction structure has exhibited broad sensitivity to illumination ranging from DUV to NIR light, with peak sensitivity to



**Table 1.** Comparison of characteristic parameters for 2D materials based photodetectors.

Device structure	$R$	$\tau_r/\tau_f$	$I_{\text{light}}/I_{\text{dark}}$	$D^*$ [Jones]	Measurement conditions	Spectral range	Ref.
PtSe <sub>2</sub> /GaAs heterojunction	262 mA W <sup>-1</sup>	5.5/6.5 μs	≈10 <sup>4</sup>	≈10 <sup>12</sup>	$V_{\text{bi}}^{\text{a)}} = 0$ V $\lambda = 808$ nm	Deep UV–NIR	This work
Few-layer BP phototransistor	4.8 mA W <sup>-1</sup>	1/4 ms	–	–	$V_{\text{ds}}^{\text{b)}} = 0.2$ V $V_{\text{g}}^{\text{c)}} = 0$ V $\lambda = 640$ nm	Visible–NIR	[3]
Graphene/GaAs nanocone array Schottky junction	1.73 mA W <sup>-1</sup>	72/122 μs	≈10 <sup>4</sup>	≈10 <sup>11</sup>	$V_{\text{bi}} = 0$ V $\lambda = 850$ nm	Visible–NIR	[18]
Multilayer MoS <sub>2</sub> phototransistor	<120 mA W <sup>-1</sup>	–	–	≈10 <sup>10</sup>	$V_{\text{ds}} = 1$ V $V_{\text{g}} = 0$ V $\lambda = 633$ nm	Visible–NIR	[38]
Graphene/ $\beta$ -Ga <sub>2</sub> O <sub>3</sub> heterojunction	39.9 A W <sup>-1</sup>	94.83/219.19 s	–	≈10 <sup>13</sup>	$V_{\text{bi}} = 2$ V $\lambda = 254$ nm	Deep UV	[35]
WS <sub>2</sub> /Si heterojunction	5.7 A W <sup>-1</sup>	670/998 μs	≈10 <sup>1</sup>	–	$V_{\text{bi}} = -5$ V $\lambda = 660$ nm	UV–NIR	[48]
MoS <sub>2</sub> /Si heterojunction	11.9 A W <sup>-1</sup>	30.5/71.6 μs	59.9	≈10 <sup>10</sup>	$V_{\text{bi}} = -2$ V $\lambda = 650$ nm	UV–NIR	[49]
Graphene/Si heterojunction	0.435 mA W <sup>-1</sup>	1.2/3 ms	≈10 <sup>4</sup>	≈10 <sup>9</sup>	$V_{\text{bi}} = 0$ V $\lambda = 850$ nm	UV–NIR	[50]
Few-layer WS <sub>2</sub> photosensor	0.092 mA W <sup>-1</sup>	5.3/5.3 ms	–	–	$V_{\text{bi}} = 5$ V $\lambda = 514$ nm	Visible	[51]
Few-layer InSe phototransistor	27 A W <sup>-1</sup>	0.5/1.7 s	<10	–	$V_{\text{ds}} = 1$ V $V_{\text{g}} = 0$ V $\lambda = 370$ nm	UV–NIR	[52]
Graphene/Ge heterojunction	51.8 mA W <sup>-1</sup>	23/108 μs	≈10 <sup>4</sup>	≈10 <sup>10</sup>	$V_{\text{bi}} = 0$ V $\lambda = 1550$ nm	NIR	[53]

<sup>a)</sup>  $V_{\text{bi}}$ : The bias voltage; <sup>b)</sup>  $V_{\text{ds}}$ : Source–drain voltage; <sup>c)</sup>  $V_{\text{g}}$ : The positive gate voltage.



**Figure 8.** a) The  $I$ – $V$  characteristics and b) zero-bias time-dependent photoresponse of the PtSe<sub>2</sub>/GaAs photodetector after long-term storage. c) Comparison of the XPS spectra of the PtSe<sub>2</sub> before and after six weeks. d) Raman spectra of the PtSe<sub>2</sub> film before and after six weeks.

illumination ranging from 650 to 810 nm at zero bias voltage. The  $I_{\text{on}}/I_{\text{off}}$  ratio, responsivity, specific detectivity, and response speed of the studied photodetector were estimated to be  $3 \times 10^4$ , 262 mA W<sup>-1</sup>,  $2.52 \times 10^{12}$  Jones, and 5.5/6.5 μs, respectively, which were comparable to or better than that of other TMDs photodetectors. Further DFT theoretical simulation results revealed that the peak sensitivity in the range from 650 nm (visible) to 810 nm (near infrared) was probably related to the interstitial Se atoms in PtSe<sub>2</sub> films, which will lead to opening the bandgap. It is expected that the high-performance PtSe<sub>2</sub>/GaAs heterojunction photodetector demonstrated in this paper will open up a new pathway for the development of next-generation 2D Group-10 materials based optoelectronic devices.

## 4. Experimental Section

**Materials Synthesis and Characterization:** The PtSe<sub>2</sub> in this work was grown by a simple selenization method. In brief, 12 nm Pt film was first deposited on SiO<sub>2</sub>/Si (300 nm SiO<sub>2</sub> thickness) using magnetron sputtering system. The metal-deposited SiO<sub>2</sub>/Si substrates were placed at the center zone of the growth furnace and elemental selenium powder (99.99% purity) were placed at the upstream side. Selenium was evaporated at 220 °C dragged by 50 SCCM (standard cubic centimeter per minute) argon flow. The center temperature of the tube furnace was set to be 420 °C. After selenization for 1 h, a thin film gray in color, will be formed at the substrate. The absorption spectra of PtSe<sub>2</sub> on quartz glass, GaAs substrate, and PtSe<sub>2</sub>/GaAs heterojunction were recorded on a Perkin Elmer Lambda 900 UV/VIS/NIR spectrometer. The Raman spectra measurements were carried out on a HORIBA Raman spectrometer with a 488 nm argon ion laser. The XRD pattern was recorded using a RigakuSmartLab X-ray diffractometer. The topography of samples was obtained by AFM (VeecoNanoscope V). The morphology of the PtSe<sub>2</sub> was observed by SEM (JEOL Model JSM-6490). The XPS measurements were performed using a monochromatic Al Kα source (1486.6 eV) produced by the XPS system. The morphology, crystal structure, and chemical composition were investigated using a field emission transmission electron microscope (FETEM, JEOL Model JEM-2100F), equipped with an energy-dispersive spectrometer.

**Device Fabrication and Characterization:** To fabricate the PtSe<sub>2</sub>/GaAs heterojunction photodetector, a 200 nm thick SiO<sub>2</sub> film was deposited on a GaAs wafer using a shadow mask to form a window (0.3 cm × 0.3 cm), in which the PtSe<sub>2</sub> film will be transferred. The PtSe<sub>2</sub> films were spin-coated with 5 wt% polymethylmethacrylate (PMMA) in chlorobenzene, and then the underlying SiO<sub>2</sub>/Si was removed in NaOH solution. The PtSe<sub>2</sub> film was rinsed in deionized water to remove the remaining ions. Afterward, the above GaAs was soaked in deionized water solution, and then slowly lifted to transfer the PtSe<sub>2</sub> films on GaAs surface. Finally, a drop of Ag paste was then placed at the centers of the PtSe<sub>2</sub> films. The optoelectronic characterization of the device was performed using a semiconductor parameter analyzer system (Keithley 4200-SCS) at room temperature. To investigate the spectral response and response speed of the heterojunction devices, a home-built optoelectronic system composed of a light source (LE-SP-LS-XE), a monochromator (LE-SP-M300), an oscilloscope (Tektronix, TDS2012B), and an optical chopper (LE-oc120) was employed.

**Computational Method:** All calculations were performed based on DFT method implemented in the Vienna ab initio simulation package (VASP).<sup>[56–59]</sup> The exchange correlation potential was approximated by the generalized gradient approximation (GGA) with Perdew–Burke–Ernzerhof (PBE) functional.<sup>[60]</sup> To describe correctly van der Waals (vdW) interaction, the DFT-D3 method with Becke–Johnson damping was employed in the calculation.<sup>[61]</sup> Electronic wave functions were expanded using a plane-wave basis set with a cutoff energy of 350 eV. The PtSe<sub>2</sub> thin films can be cleaved from the (0001) surface of the bulk which

belongs to the  $D_{3d}^3$  ( $P\bar{3}m1$ ) space group of the trigonal system. In the PtSe<sub>2</sub> thin films with one layer, one atomic layer of Pt was sandwiched between two atomic layers of Se. For the PtSe<sub>2</sub> thin films with one, two, and three layers, the  $1 \times 1$  primitive cells were chosen and a  $k$ -point set of  $40 \times 40 \times 1$  was employed in geometrical optimization and electronic structure calculations. For the PtSe<sub>2</sub> thin films with three layers, the  $2 \times 2$ ,  $3 \times 3$ , and  $4 \times 4$  supercells were used to investigate the doping effects on the electronic structures. In geometrical and electronic structure calculations, the  $k$ -point sets of  $35 \times 35 \times 1$ ,  $25 \times 25 \times 1$ , and  $15 \times 15 \times 1$  were employed for these three supercells, respectively. A vacuum space of 20 Å was used along the  $z$ -direction perpendicular to 2D planes to eliminate the interaction among the periodic images. The full relaxation of structures was carried out until the force component on each atom is less than 0.01 eV Å<sup>-1</sup>, and the convergence criterion is 10<sup>-5</sup> eV for energy.

## Supporting Information

Supporting Information is available from the Wiley Online Library or from the author.

## Acknowledgements

L.-H.Z. and S.-H.L. contributed equally to this work. This work was supported by the National Natural Science Foundation of China (NSFC, Nos. 61675062, 61575059, and 21501038), the Fundamental Research Funds for the Central Universities (2013HGCH0012, 2014HGCH0005), the China Postdoctoral Science Foundation (103471013), the Research Grants Council of Hong Kong, China (Project No. GRF 152109/16E PolyU B-Q52T), and the Hong Kong Polytechnic University (Project Nos. 4-BCCW and 1-ZVGH).

## Conflict of Interest

The authors declare no conflict of interest.

## Keywords

broadband, density functional theory, heterojunctions, photodetectors, transitional metal dichalcogenides

Received: October 15, 2017

Revised: January 4, 2018

Published online:

- [1] X. Yin, Z. Ye, D. A. Chenet, Y. Ye, K. O'Brien, J. C. Hone, X. Zhang, *Science* **2014**, *344*, 488.
- [2] P. Wang, S. Liu, W. Luo, H. Fang, F. Gong, N. Guo, Z. G. Chen, J. Zou, Y. Huang, X. Zhou, J. Wang, X. Chen, W. Lu, F. Xiu, W. Hu, *Adv. Mater.* **2017**, *29*, 1.
- [3] M. Buscema, D. J. Groenendijk, S. I. Blanter, G. A. Steele, H. S. van der Zant, A. Castellanos-Gomez, *Nano Lett.* **2014**, *14*, 3347.
- [4] C. Yan, L. Gan, X. Zhou, J. Guo, W. Huang, J. Huang, B. Jin, J. Xiong, T. Zhai, Y. Li, *Adv. Funct. Mater.* **2017**, *27*, 1702918.
- [5] F.-X. Liang, L.-B. Luo, *Sci. China-Phys. Mech. Astron.* **2016**, *60*, 037031.
- [6] C. Xie, C. Mak, X. Tao, F. Yan, *Adv. Funct. Mater.* **2017**, *27*, 1.

- [7] J. Chu, F. Wang, L. Yin, L. Lei, C. Yan, F. Wang, Y. Wen, Z. Wang, C. Jiang, L. Feng, J. Xiong, Y. Li, J. He, *Adv. Funct. Mater.* **2017**, *27*, 1701342.
- [8] X. Chia, A. Adriano, P. Lazar, Z. Sofer, J. Luxa, M. Pumera, *Adv. Funct. Mater.* **2016**, *26*, 4306.
- [9] M. Zhang, R. C. T. Howe, R. I. Woodward, E. J. R. Kelleher, F. Torrisi, G. Hu, S. V. Popov, J. R. Taylor, T. Hasan, *Nano Res.* **2015**, *8*, 1522.
- [10] K. C. Kwon, S. Choi, K. Hong, C. W. Moon, Y.-S. Shim, D. H. Kim, T. Kim, W. Sohn, J.-M. Jeon, C.-H. Lee, K. T. Nam, S. Han, S. Y. Kim, H. W. Jang, *Energy Environ. Sci.* **2016**, *9*, 2240.
- [11] Y. Zhao, J. Qiao, Z. Yu, P. Yu, K. Xu, S. P. Lau, W. Zhou, Z. Liu, X. Wang, W. Ji, Y. Chai, *Adv. Mater.* **2017**, *29*, 1.
- [12] Y. Wang, L. Li, W. Yao, S. Song, J. T. Sun, J. Pan, X. Ren, C. Li, E. Okunishi, Y. Q. Wang, E. Wang, Y. Shao, Y. Y. Zhang, H. T. Yang, E. F. Schwier, H. Iwasawa, K. Shimada, M. Taniguchi, Z. Cheng, S. Zhou, S. Du, S. J. Pennycook, S. T. Pantelides, H. J. Gao, *Nano Lett.* **2015**, *15*, 4013.
- [13] C. Yim, K. Lee, N. McEvoy, M. O'Brien, S. Riazimehr, N. C. Berner, C. P. Cullen, J. Kotakoski, J. C. Meyer, M. C. Lemme, *ACS Nano* **2016**, *10*, 9550.
- [14] A. Ali Umar, S. K. Md Saad, M. Mat Salleh, *ACS Omega* **2017**, *2*, 3325.
- [15] C. Yim, V. Passi, M. C. Lemme, G. S. Duesberg, C. Ó. Pallechi, D. Fadil, N. McEvoy, arXiv preprint arXiv:1707.06824, **2017**.
- [16] Z. G. Wang, Q. Li, F. Besenbacher, a. M. D. Dong, *Adv. Mater.* **2016**, *28*, 10224.
- [17] K. Ullah, S. Ye, Z. Lei, K.-Y. Cho, W.-C. Oh, *Catal. Sci. Technol.* **2015**, *5*, 184.
- [18] L.-B. Luo, J.-J. Chen, M.-Z. Wang, H. Hu, C.-Y. Wu, Q. Li, L. Wang, J.-A. Huang, F.-X. Liang, *Adv. Funct. Mater.* **2014**, *24*, 2794.
- [19] L. B. Luo, L. H. Zeng, C. Xie, Y. Q. Yu, F. X. Liang, C. Y. Wu, L. Wang, J. G. Hu, *Sci. Rep.* **2014**, *4*, 3914.
- [20] L. Tang, R. Ji, X. Li, G. Bai, C. P. Liu, J. Hao, J. Lin, H. Jiang, K. S. Teng, Z. Yang, S. P. Lau, *ACS Nano* **2014**, *8*, 6312.
- [21] P. Norton, *Opto-Electron. Rev.* **2002**, *10*, 159.
- [22] R. Sporken, S. Sivananthan, K. K. Mahavadi, G. Monfroy, M. Boukerche, J. P. Faurie, *Appl. Phys. Lett.* **1989**, *55*, 1879.
- [23] L. Ye, H. Li, Z. Chen, J. Xu, *ACS Photonics* **2016**, *3*, 692.
- [24] V. M. Bazovkin, S. A. Dvoretzkiy, A. A. Guzev, A. P. Kovchavtsev, D. V. Marin, Z. V. Panova, I. V. Sabinina, Y. G. Sidorov, G. Y. Sidorov, A. V. Tsarenko, V. S. Varavin, V. V. Vasiliev, M. V. Yakushev, *Phys. Status Solidi C* **2016**, *13*, 651.
- [25] S. Ghosh, S. M. Han, *IEEE Electron Dev. Lett.* **2014**, *35*, 900.
- [26] L. Tao, H. Long, B. Zhou, S. F. Yu, S. P. Lau, Y. Chai, K. H. Fung, Y. H. Tsang, J. Yao, D. Xu, *Nanoscale* **2014**, *6*, 9713.
- [27] M. O'Brien, N. McEvoy, C. Motta, J.-Y. Zheng, N. C. Berner, J. Kotakoski, K. Elibol, T. J. Pennycook, J. C. Meyer, C. Yim, M. Abid, T. Hallam, J. F. Donegan, S. Sanvito, G. S. Duesberg, *2D Mater.* **2016**, *3*, 021004.
- [28] Y. Jung, J. Shen, Y. Liu, J. M. Woods, Y. Sun, J. J. Cha, *Nano Lett.* **2014**, *14*, 6842.
- [29] D. Kong, H. Wang, J. J. Cha, M. Pasta, K. J. Koski, J. Yao, Y. Cui, *Nano Lett.* **2013**, *13*, 1341.
- [30] W.-C. Oh, K. Ullah, L. Zhu, Z.-D. Meng, S. Ye, S. Sarkar, *Mater. Sci. Semicond. Process.* **2014**, *25*, 34.
- [31] L.-B. Luo, H. Hu, X.-H. Wang, R. Lu, Y.-F. Zou, Y.-Q. Yu, F.-X. Liang, *J. Mater. Chem. C* **2015**, *3*, 4723.
- [32] W. Jie, F. Zheng, J. Hao, *Appl. Phys. Lett.* **2013**, *103*, 233111.
- [33] H. Sun, T. Lei, W. Tian, F. Cao, J. Xiong, L. Li, *Small* **2017**, *13*, 1701042.
- [34] G. Konstantatos, I. Howard, A. Fischer, S. Hoogland, J. Clifford, E. Klem, L. Levina, E. H. Sargent, *Nature* **2006**, *442*, 180.
- [35] W. Y. Kong, G. A. Wu, K. Y. Wang, T. F. Zhang, Y. F. Zou, D. D. Wang, L. B. Luo, *Adv. Mater.* **2016**, *28*, 10725.
- [36] H. Kind, H. Yan, B. Messer, M. Law, P. Yang, *Adv. Mater.* **2002**, *14*, 158.
- [37] Y. Zhang, Y. Yu, X. Wang, G. Tong, L. Mi, Z. Zhu, X. Geng, Y. Jiang, *J. Mater. Chem. C* **2017**, *5*, 140.
- [38] W. Choi, M. Y. Cho, A. Konar, J. H. Lee, G. B. Cha, S. C. Hong, S. Kim, J. Kim, D. Jena, J. Joo, S. Kim, *Adv. Mater.* **2012**, *24*, 5832.
- [39] J. Mao, Y. Yu, L. Wang, X. Zhang, Y. Wang, Z. Shao, J. Jie, *Adv. Sci.* **2016**, *3*, 1600018.
- [40] D. Melisi, M. A. Nitti, M. Valentini, A. Valentini, T. Ligonzo, G. De Pascali, M. Ambrico, *Beilstein J. Nanotechnol.* **2014**, *5*, 1999.
- [41] P. Li, L. Li, X. C. Zeng, *J. Mater. Chem. C* **2016**, *4*, 3106.
- [42] H. B. Zhang, X. J. Zhang, C. Liu, S. T. Lee, J. S. Jie, *ACS Nano* **2016**, *10*, 5113.
- [43] Y. Park, V. Choong, Y. Gao, B. R. Hsieh, C. W. Tang, *Appl. Phys. Lett.* **1996**, *68*, 2699.
- [44] L. Wang, J. Jie, Z. Shao, Q. Zhang, X. Zhang, Y. Wang, Z. Sun, S.-T. Lee, *Adv. Funct. Mater.* **2015**, *25*, 2910.
- [45] L. Zeng, C. Xie, L. Tao, H. Long, C. Tang, Y. H. Tsang, J. Jie, *Opt. Express* **2015**, *23*, 4839.
- [46] B. Nie, J. G. Hu, L. B. Luo, C. Xie, L. H. Zeng, P. Lv, F. Z. Li, J. S. Jie, M. Feng, C. Y. Wu, Y. Q. Yu, S. H. Yu, *Small* **2013**, *9*, 2872.
- [47] A. Osinsky, S. Gangopadhyay, R. Gaska, B. Williams, M. A. Khan, D. Kuksenkov, H. Temkin, *Appl. Phys. Lett.* **1997**, *71*, 2334.
- [48] C. Lan, C. Li, S. Wang, T. He, T. Jiao, D. Wei, W. Jing, L. Li, Y. Liu, *ACS Appl. Mater. Interfaces* **2016**, *8*, 18375.
- [49] Y. Zhang, Y. Yu, L. Mi, H. Wang, Z. Zhu, Q. Wu, Y. Zhang, Y. Jiang, *Small* **2016**, *12*, 1062.
- [50] X. An, F. Liu, Y. J. Jung, S. Kar, *Nano Lett.* **2013**, *13*, 909.
- [51] N. Perea-López, A. L. Elías, A. Berkdemir, A. Castro-Beltran, H. R. Gutiérrez, S. Feng, R. Lv, T. Hayashi, F. López-Urías, S. Ghosh, B. Muchharla, S. Talapatra, H. Terrones, M. Terrones, *Adv. Funct. Mater.* **2013**, *23*, 5511.
- [52] Z. Yang, W. Jie, C. H. Mak, S. Lin, H. Lin, X. Yang, F. Yan, S. P. Lau, J. Hao, *ACS Nano* **2017**, *11*, 4225.
- [53] L. H. Zeng, M. Z. Wang, H. Hu, B. Nie, Y. Q. Yu, C. Y. Wu, L. Wang, J. G. Hu, C. Xie, F. X. Liang, L. B. Luo, *ACS Appl. Mater. Interfaces* **2013**, *5*, 9362.
- [54] C. Xie, B. Nie, L. Zeng, F.-X. Liang, M.-Z. Wang, L. Luo, M. Feng, Y. Yu, C.-Y. Wu, Y. Wu, *ACS Nano* **2014**, *8*, 4015.
- [55] J. D. Wood, S. A. Wells, D. Jariwala, K. S. Chen, E. Cho, V. K. Sangwan, X. Liu, L. J. Lauhon, T. J. Marks, M. C. Hersam, *Nano Lett.* **2014**, *14*, 6964.
- [56] P. E. Blöchl, *Phys. Rev. B* **1994**, *50*, 17953.
- [57] G. Kresse, J. Furthmüller, *Comput. Mater. Sci.* **1996**, *6*, 15.
- [58] G. Kresse, D. Joubert, *Phys. Rev. B* **1999**, *59*, 1758.
- [59] Z. Li, W. Xu, Y. Yu, H. Du, K. Zhen, J. Wang, L. Luo, H. Qiu, X. Yang, *J. Mater. Chem. C* **2016**, *4*, 362.
- [60] J. P. Perdew, K. Burke, M. Ernzerhof, *Phys. Rev. Lett.* **1996**, *77*, 3865.
- [61] S. Grimme, J. Antony, S. Ehrlich, H. Krieg, *J. Chem. Phys.* **2010**, *132*, 154104.