# **RSC Advances**

## PAPER



Cite this: RSC Adv., 2015, 5, 13324

p-type ZnTe:Ga nanowires: controlled doping and optoelectronic device application

CHEMISTRY

View Article Online

View Journal | View Issue

Lin-Bao Luo,<sup>\*a</sup> Shun-Hang Zhang,<sup>a</sup> Rui Lu,<sup>a</sup> Wei Sun,<sup>b</sup> Qun-Ling Fang,<sup>c</sup> Chun-Yan Wu,<sup>\*a</sup> Ji-Gang Hu<sup>a</sup> and Li Wang<sup>\*a</sup>

Although significant progress has been achieved in the synthesis and doping of ZnTe nanostructures, it remains a major challenge to rationally tune their transport properties for nanodevice applications. In this work, p-type ZnTe nanowires (NWs) with tunable conductivity were synthesized by employing Ga/Ga<sub>2</sub>O<sub>3</sub> as a dopant *via* a simple thermal evaporation method. Electrical measurements of back-gate metal-oxide field-effect-transistors based on a single NW revealed that when the Ga content in the ZnTe NWs increases from 1.3 to 5.1 and 8.7%, the hole mobility and hole concentration will increase from 0.0069 to 0.33 to 0.46 cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup>, respectively. It was also found that the photodetector composed of a ZnTe:Ga NW/graphene Schottky diode exhibited high sensitivity to visible light illumination with an on/off ratio as high as 10<sup>2</sup> at reverse bias, with good reproducibility. The responsivity and detectivity were estimated to be 4.17 × 10<sup>3</sup> A W<sup>-1</sup> and 3.19 × 10<sup>13</sup> cm Hz<sup>1/2</sup> W<sup>-1</sup>, higher than other ZnTe nanostructure based photodetectors. It is expected that the ZnTe:Ga NWs with controlled p-type conductivity are promising building blocks for fabricating high performance nano-optoelectronic devices in the future.

Received 8th November 2014 Accepted 19th January 2015 DOI: 10.1039/c4ra14096f

www.rsc.org/advances

#### Introduction

Over the past decade, one dimensional (1-D) semiconductor nanostructures with high aspect ratio, such as nanowires (NWs), nanoribbons (NRs), nanotubes (NTs), and multibranched structures, have attracted considerable research interest because of their novel physical and chemical properties. Among the huge semiconductor nanostructures family, II-VI group semiconductor nanomaterials (e.g., ZnS, ZnSe, CdS, CdSe, CdTe) have received special attention, and their unique optical properties have intensively investigated.<sup>1-3</sup> It has been reported that one-dimensional II-VI group nanostructures are good building blocks for fabricating high performance nanoelectronic and nano-optoelectronic devices,4,5 including nano field-effect transistors (nano-FETs),6,7 nano photodetectors (nano-PDs),<sup>8-10</sup> nano light emitting diodes (nano-LEDs),<sup>11,12</sup> non-volatile memory device,13,14 nanogenerator,15 solar cells and so on.16,17

Zinc telluride (ZnTe), as an important II–VI group semiconductor, has a moderate and direct band-gap of 2.26 eV at room temperature<sup>18-20</sup> due to improved electrical and optical

properties, ZnTe nanostructures in one-dimensional form have demonstrated great potential for a variety of optoelectronic devices,<sup>21</sup> including green light-emitting diodes,<sup>22</sup> electro-optic detectors,<sup>23</sup> and solar cells.<sup>24</sup> To date, a number of methods have been developed to tune the conductivity, such as electrochemical deposition,25 molecular beam epitaxy,26 thermal evaporation,27 and metal organic chemical vapor deposition.28 In order to increase the p-type conductivity, various impurity atoms including Cu, P, Sb, and N have been introduced into the crystal lattice of ZnTe during growth process.<sup>10,29</sup> For instance, Li et al. reported N doped p-type ZnTe nanoribbons via a thermal evaporation in a mixed atmosphere of hydrogen, ammonia and argon. Electrical analysis revealed that the nitrogen doping led to a substantial enhancement in hole mobility from 6.2  $\times$  10<sup>-3</sup> to 1.2 cm V s<sup>-1</sup>.<sup>30</sup> Recently, we developed a thermal evaporation method to achieve Sb doped ZnTe NWs by using antimony solid source. Moreover, the p-type conductivity of ZnTe NWs can be also enhanced by dipping the intrinsic ZnTe NWs with  $Cu(NO_3)_2$  solution, followed by a post-annealing process to allow the diffusion of the Cu ions into the ZnTe NWs.<sup>31</sup> In spite of these progress, unfortunately, the use of environmentally unfriendly Sb, toxic ammonia (NH<sub>3</sub>) gas inevitably increases the complexity of the synthesis process, and limits their applications. On the other hand, the Cu ions solution method lacks of reproducibility and controllability. Therefore, a simple, efficient and environmental friendly method to dope ZnTe nanostructures with well-controlled ptype conductivity is highly desired. Herein, we present the synthesis of ZnTe NWs with enhanced p-type conductivity via a

<sup>&</sup>quot;School of Electronic Sciences and Applied Physics and Anhui Provincial Key Laboratory of Advanced Functional Materials and Devices, Hefei University of Technology, Hefei, Anhui 230009, P. R. China. E-mail: luolb@hfut.edu.cn; wlhgd@ hfut.edu.cn; cywu@hfut.edu.cn

<sup>&</sup>lt;sup>b</sup>School of Materials Science and Engineering, Hefei University of Technology, Hefei, Anhui 230009, P. R. China

<sup>&</sup>lt;sup>c</sup>School of Medical Engineering, Hefei University of Technology, Hefei, Anhui 230009, P. R. China

#### Paper

simple thermal evaporation method by using Ga/Ga<sub>2</sub>O<sub>3</sub> as dopant. Electrical measurement of the single ZnTe:Ga NW revealed that the conductivities of the NWs could be tuned in a wide range by adjusting the doping concentration. In addition, nano-PD based on ZnTe:Ga NW/graphene Schottky diode exhibited high sensitivity to visible light illumination with good reproducibility.

### **Experimental details**

The fabrication of Ga-doped ZnTe NWs was conducted in a horizontal tube furnace via a thermal evaporation method, in which ZnTe and a mixed powder of Ga and  $Ga_2O_3$  (4 : 1 molar ratio) were used as source and dopant, respectively. Briefly, an alumina boat containing 0.5 g ZnTe powder mixed with dopant was firstly placed into the central position of the furnace. After being evacuated to a base pressure of  $10^{-3}$  Torr, the reaction tube was flushed and filled with a mixture gas of 20 standardstate cubic centimeter per minute (Sccm) Ar and H<sub>2</sub> (5% in volume). When the pressure in reaction chamber was kept at 160 Torr, the Ga/Ga<sub>2</sub>O<sub>3</sub> and the ZnTe sources were then heated up to 1050 °C, and maintained at the temperature for 2 h. After the system was cooled down to room temperature, the Si substrates with a lays of wool-like product was taken out from the furnace. During the growth process, the Ga doping levels can be easily controlled by tailoring the Ga/Ga<sub>2</sub>O<sub>3</sub> ratio in the sources. In this study, totally three samples with varied Ga doping levels were fabricated and marked as sample A, B and C, corresponding to the increasing Ga:Ga2O3 powder mixture to ZnTe mass ratio of 2%, 4% and 8%, respectively.

The morphologies and micro-structures of the ZnTe:Ga NWs were characterized by X-ray diffraction (XRD, Rigaku D/Max 2500V), scanning electron microscopy (SEM, JSM-6490LV) and high-resolution transmission electron microscopy (HRTEM, JEOL JEM-2100F) with selected area electron diffraction (SAED). The chemical composition of the as-synthesized NWs was determined by energy-dispersive X-ray spectroscopy (EDX, attached on the SEM) and X-ray photoelectron spectroscopy (XPS, Thermo ESCALAB250) which was equipped with a monochromatic Al X-ray (1486.6 eV) source.

To evaluate the effect of Ga doping on the electrical properties of ZnTe NWs, back-gate field-effect transistors (FETs) based on single NW were fabricated. Firstly, the above ZnTe:Ga NWs were dispersed on SiO<sub>2</sub> (300 nm)/p<sup>+</sup>-Si substrate with a desired density. Then photolithography and electron beam evaporation processes were used to define the source and drain electrodes. Finally, the remaining photoresist was removed through a simple lift-off process. Besides the above FETs device, Schottky diodes based on ZnTe:Ga NW/graphene heterojunctions were fabricated as well. The as-synthesized ZnTe:Ga NWs were parallel dispersed on the substrate by contact print technique, followed by fabrication of single Cu/Au (50 nm) electrode on the ZnTe:Ga NW via photolithography, electron beam evaporation and lift-off process. Afterwards, monolayer graphene was transferred onto the one head of the NW. In order to form a good contact to the graphene, silver paste was finally placed onto the graphene. All the electrical measurements on the nano-FETs and Schottky diodes were conducted at room temperature by using a semiconductor characterization system (Keithley 4200-SCS).

#### **Results and discussion**

Fig. 1(a) shows the digital camera picture of both intrinsic and Ga doped ZnTe NWs, according to which all samples appear nearly the same in color, suggesting that the Ga doping can hardly alter the color of ZnTe NWs. Fig. 1(b) displays a typical SEM image of intrinsic ZnTe NWs, it is obvious that the product is mainly composed of straight fiber-like nanostructures, with diameters of 200–400 nm, and lengths of up to several tens of micrometers. The Zn/Te atomic ratio of these ZnTe NWs is estimated to be about 48 : 52 (the inset of Fig. 1(b)), which is close to the stoichiometric ratio of ZnTe. Remarkably, after Ga doping, all the samples (sample A, B, and C) nearly retain the morphology of the NWs.

According to the XRD patterns of both intrinsic and the Gadoped ZnTe NWs (sample A, B and C) in Fig. 2(a), all the peaks can be readily assigned to ZnTe (JCPDS no. 15-0746), and no obvious peak shift is observed, suggesting that the product is of high quality without contaminant phase. Interestingly, care examination of the (111) signal observed obvious shift to large 2 theta value when the doping level was gradually increased (Fig. 2(b)). Such shift is understandably due to the relatively small Ga ions compared with Zn ions. When Ga ions with small diameter substituted the Zn site in the ZnS crystal, the resultant crystal unit cell become smaller, thus the 2 theta is increased according to Bragg's Law. To study the Ga concentration in the ZnTe NWs, X-ray photoelectron spectroscopy (XPS) spectra of all produces were studied (Fig. 2(c-e)). Obviously, after Ga doping a



Fig. 1 (a) Digital camera pictures of both intrinsic and Ga doped ZnTe NWs. (b) SEM images of intrinsic ZnTe NWs, the inset shows the corresponding EDX spectrum. (c)–(e) represent ZnTe NWs with different Ga doping levels: (c) sample A, (d) sample B and (e) sample C.



Fig. 2 (a) XRD patterns of both intrinsic and Ga doped ZnTe NWs. (b) Enlarged XRD pattern of the (111). (c)–(e) are XPS spectra of sample A, sample B, and sample C, respectively, the insets show the corresponding Ga 3d peak.

weak peak at  $\sim$ 20 eV attributable to the Ga 3d core level emission is observed. Careful examination of these peaks found that the Ga atoms were doped into ZnTe NWs in the form of +3 valence. Moreover, the Ga contents in the sample A, B and C are estimated to be 1.3%, 5.1%, and 8.7% (at%), respectively, from the XPS spectrum (the insetting figures).

Fig. 3(a) shows a TEM image of a single NW from sample A, with a diameter of around 150 nm. The smooth surface signifies good crystallinity and uniformity of Ga-doped ZnTe NWs. Further lattice resolved image of HRTEM in Fig. 3(b) displays a single-crystal zinc blende phase, the *d*-spacing value of the plane perpendicular to the growth direction is estimated to be of 0.36 nm, corresponding to the [200] direction. As matter of fact, such a growth direction was further confirmed by the selected election diffraction (SAED) pattern shown in the inset of Fig. 3(b). To unveil how the impurity atoms were distributed in the ZnTe NW, we studied the elemental mapping of Zn, Te,



**Fig. 3** (a) Typical TEM image of an individual ZnTe NW from sample A. (b) HRTEM image of the ZnTe:Ga NW, the inset shows the corresponding SAED pattern. (c–e) Elemental mapping of an individual NW from sample A, respectively.

and Ga in a typical ZnTe from sample A. As illustrated in Fig. 3(c-e), all the constituting elements are uniformly distributed in the NW. Notably, due to the relatively low concentration, the mapping of Ga atoms is much vaguer, compared with that of Zn and Te.

To explore the Ga doping dependent transport properties of the NWs, nano-FET based on individual ZnTe:Ga NW was fabricated (the insets of Fig. 4(a)). Fig. 4(a) plots the typical I-V



**Fig. 4** (a) *I–V* curves of ZnTe:Ga NWs with varied doping levels, the above left inset shows a representative SEM image of a single ZnTe:Ga NW connected by two parallel Cu/Au electrodes, the below right inset is the schematic illustration of the back-gate FET device based on a single ZnTe:Ga NW. (b) Conductivity distribution of the ZnTe:Ga NWs with different doping levels.

curves measured from the samples with different doping levels. The ideal linearity of the *I*-*V* curves confirms the ohmic contact formed between the ZnTe:Ga NWs and the Cu/Au electrodes. The conductivity of the ZnTe NWs from Sample A (doping percentage: 1.3%) was  $2.7 \times 10^{-4}$  S cm<sup>-1</sup>, which is actually a little bit smaller than that without doping  $(0.0016 \text{ S cm}^{-1})$ .<sup>25</sup> When the doping concentration was further increased to 5.1%, and 8.7%, the corresponding conductivities in contrast begin to increase to 0.069 S cm<sup>-1</sup> and 0.43 S cm<sup>-1</sup>, respectively. We attribute this relatively small conductivity of sample A to the weak self-compensation effect at low doping level: as observed on one-dimensional II-VI semiconductor nanostructures,32 one can easily achieve n-type doping for CdS, ZnO and CdSe, and p-type doping for ZnTe, but it is nearly impossible to realize p-type doping for CdS, ZnO and CdSe, and n-type doping for ZnTe. This difficulty mainly arises from the strong self-compensation effect, that is, doping of the acceptors will inevitably introduce one or more acceptor (donor) defects in the lattice, thus counteracting the effect of n-type doping (p-type). In this study, for NWs with low Ga doping levels, the self-compensation effect is too weak to generate sufficient holes, thus the conductivity is even higher than the intrinsic NWs. To gain a statistical significance, the electrical conductivities of totally 15 NWs were measured. As shown in Fig. 4(b), most of the NWs from sample A showed conductivity in the magnitude of 10<sup>-4</sup> S cm<sup>-1</sup>. For sample B, most NWs have a conductivity value in the range of 1–10 S  $\rm cm^{-2}$ , and for sample C, it increases to the range of 0.1–1 S cm $^{-1}$ .

Fig. 5 shows the representative  $I_{DS}-V_{DS}$  curves of the ZnTe:Ga NWs measured at varied gate voltages ( $V_G$ ) ranging from -40 to



**Fig. 5** Electron-transfer characteristics of the back-gate FETs based on ZnTe:Ga NWs with different doping levels. The  $I_{DS}-V_{DS}$  curves were measured under varied  $V_{G}$  values ranging from -40 V to +40 V with a voltage step of 20 V. (b), (d) and (f) the  $I_{DS}-V_{G}$  curves of three samples were measured at fixed  $V_{DS} = 5$  V.

+40 V at a step of 20 V, along with the transfer characteristic at  $V_{\rm DS}$  of 5 V. For all the samples, the conductance of the NWs consistently increases (decreases) with the decreasing (increasing) of  $V_{\rm G}$ , indicative of typical p-type feature of the Gadoped ZnTe NWs. As mentioned above, the enhanced p-type conductivity for the ZnTe:Ga NWs should be attributed to the strong compensation effect which will induce more holes when Ga atoms are incorporated into the crystal lattice. Furthermore, the hole mobility ( $\mu_{\rm h}$ ) of the ZnTe:Ga NWs can be estimated from the channel transconductance ( $g_{\rm m}$ ) of the nano-FETs ( $g_{\rm m} = dI_{\rm DS}/dV_{\rm GS}$ ) by the following equation:

$$u_{\rm h} = g_{\rm m} \frac{\ln(4h/d)L}{2\pi\varepsilon_0\varepsilon_{{\rm SiO}_2}V_{\rm DS}}$$

where  $\varepsilon$  is the dielectric constant of the gate SiO<sub>2</sub> (3.9),  $\varepsilon_0$  is vacuum permittivity. *L*, *h* and *d* represent the NW channel length (15 µm), the gate oxide layer thickness (300 nm), and the NWs diameters (~320 nm for sample A, ~200 nm for sample B and ~400 nm for sample C), respectively. Based on these values,  $\mu_h$  are estimated to be 0.0069, 0.33 and 0.46 cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup> for sample A, B and C, respectively. This doping level dependent hole mobility agrees well with what was observed in doping of ZnSe NWs. What is more, hole concentrations (*h*) in the NWs are estimated to be 2.4 × 10<sup>17</sup>, 1.3 × 10<sup>18</sup>, and 5.6 × 10<sup>18</sup> cm<sup>-3</sup> for sample A, sample B and sample C, respectively, by using the following equation:

$$n = \frac{\sigma}{\mu_{\rm h} q}$$

where  $\sigma$  is the NW's conductivity, q is the elementary charge, and  $\mu_{\rm h}$  is the hole mobility. Table 1 summarizes the electrical properties of ZnTe:Ga NW with varied doping levels. One can clearly see that with the increase of Ga content in ZnTe NWs, the transconductance, conductivity, hole mobility and concentration are all improved considerably. These results suggest that Ga could be an efficient p-type dopant for ZnTe nanostructures.

Fig. 6 exhibits the time response of the ZnTe:Ga NWs of sample A, B and C when the light illumination was switched on and off repeatedly at an external bias voltage of 5 V. It can be observed that the photoconductive characteristics of the ZnTe:Ga NWs is highly dependent on the doping levels of Ga: for sample A, the NW is sensitive to light illumination with an on/off ratio of ~10. To study the substrate surface effect on the photoresponse, we also investigate the Cu/Au/SiO<sub>2</sub>/Si junction. As shown in Fig. 6, no photoresponse was observed when a light was irradiated, suggesting that the above junction cannot contribute to the photoresponse.<sup>33,34</sup> When the doping level is slightly increased from 1.3 to 5.1, the photoresponse of the

 Table 1
 Key device parameters of the back-gate FETs based on

 ZnTe:Ga NWs with different doping levels

| Materials                        | $g_{\rm m}$ (nS)      | $\mu_{\rm h} \left( {\rm cm}^2 ~{\rm V}~{\rm s}^{-1} \right)$ | $ ho \left( \Omega 	imes \mathrm{cm}  ight)$ | $n ({\rm cm}^{-3})$  | $V_{\rm th}$ (V) |
|----------------------------------|-----------------------|---|--|--|------------------|
| Sample A<br>Sample B<br>Sample C | 0.038<br>1.37<br>3.05 | 0.0069<br>0.33<br>0.46  | $3.64 	imes 10^3$<br>14.5<br>2.42            | $egin{array}{c} 2.4 	imes 10^{17} \ 1.3 	imes 10^{18} \ 5.6 	imes 10^{18} \end{array}$ | $-10\\27\\28$    |



Fig. 6 Photoresponse of the Au electrode on SiO<sub>2</sub>/Si wafer, the intrinsic ZnTe and ZnTe:Ga NWs with varied doping levels at  $V_{DS} = 5$  V.

device is substantially reduced, and the corresponding on/off ratio is  $\sim$ 5. With further increase in doping level (8.7%), the as-fabricated device is virtually blind to light illumination. Understandably, such blindness for high doping concentration can be ascribed to the high dark current at high doping level, which will probably cover the photocurrent.<sup>35</sup>

In order to explore the potential for optoelectronic device applications, monolayer graphene and p-type ZnTe:Ga NWs from sample A were used to construct a Schottky diodes photodetector. Fig. 7(a) illustrates the schematic device structure of the ZnTe:Ga NW/graphene Schottky diodes. One can see that the ZnTe:Ga NW forms a Schottky junction with the graphene layer on one hand, and forms Ohimic contact to the Cu/ Au electrodes on the other hand (inset of Fig. 7(b)), Fig. 7(b) depicts the *I–V* characteristics of the ZnTe:Ga NW/graphene Schottky diodes measured in the dark and under illumination of both white light and 532 nm. It is obvious that at both forward and reverse bias voltages, the photocurrent will increase substantially upon light irradiation. Notably, the  $I_{\text{light}}/I_{\text{dark}}$  ratio at reserves bias voltage is as high as  $10^2$ , much higher than that at forward bias (~4). This relatively large ratio is obviously due to the low leakage current of the ZnTe:Ga NW/ graphene Schottky junction at reverse voltage. In addition to the high sensitivity, the present device also exhibit excellent stability and reproducibility at both forward and reverse bias voltages (Fig. 7(c–e), suggesting the potential application for future optoelectronic devices.

In order to quantitatively evaluate the device performance of the ZnTe:Ga NW/graphene Schottky photodetector, both responsivity R (A W<sup>-1</sup>) and detectivity  $D^*$  (cm Hz<sup>1/2</sup> W<sup>-1</sup>) were calculated by using the following equations:

$$R(\mathbf{A} \mathbf{W}^{-1}) = \left(\frac{I_{\mathrm{p}}}{P_{\mathrm{opt}}}\right) = \eta\left(\frac{q\lambda}{hc}\right)G$$

$$D^* = A^{1/2} R / (2qI_d)^{1/2} = A^{1/2} (I_p / I_{opt}) / (2qI_d)^{1/2}$$

where  $I_{\rm p}$  is the photocurrent,  $\eta$  is the quantum efficiency, h is the Planck's constant and  $\lambda$  is the light wavelength. A, q and  $I_{\rm d}$ are the PD area, elementary charge and the dark current, respectively. R and  $D^*$  are estimated to be =  $4.87 \times 10^5$  A W<sup>-1</sup> and  $3.19 \times 10^{13}$  cm Hz<sup>1/2</sup> W<sup>-1</sup>, respectively. These values are much higher than the devices composed of pure ZnTe NW,<sup>28</sup> ZnTe:Sb nanoribbons/Si nano-heterojunction,<sup>27</sup> and ZnTe nanoribbons/CdSe nano-heterojunction.<sup>36</sup>

To interpret the photoresponse characteristics of the ZnTe:Ga NW/graphene Schottky diodes, we drew the energy band diagrams of the diode at different conditions. As illustrated in Fig. 8, when graphene was coated on the p-ZnTe NW, the band of the p-ZnTe NW near the metal/semiconductor interface is bended downwards and the holes are depleted in the near-surface area of the NW. At forward bias, the Fermi level of p-ZnTe NW moves downwards, allowing the holes in the NW to pass through the Schottky barrier and reach the electrode,



Fig. 7 (a) Schematic illustration of the ZnTe:Ga NW/graphene Schottky diode. (b) I-V characteristics of the device with and without light illumination, the inset shows a representative SEM image of a single ZnTe:Ga NW/monolayer graphene Schottky diode. Time response of the device illuminated by white light at forward bias (c), and under reverse bias (d). Time response of the device illuminated by 532 nm light at forward bias (e), and under reverse bias (d).



Fig. 8 Energy band diagram of ZnTe NW/graphene Schottky diodes at (a) zero bias, (b) forward bias, (c) reverse bias.

which means "On" state for the diode. When the light illumination is switched on, the NW will absorb the photons incident on the surface, leading to the generation of electron-hole pairs, which will then be separated by the electric filed between the graphene and ZnTe:Ga NW and form the photocurrent. At reverse bias, the energy band of p-ZnTe:Ga NW will bend downwards and a larger space-charge region is formed. The electric field in the space charge region has the direction from the graphene to the semiconductor (p-ZnTe:Ga NW). Due to the large Schottky barrier at the interface, it is very difficult for holes to drift from the graphene to the NW, resulting in a low dark current. Upon light illumination, the photogenerated electronhole pairs will be separated in the space-charge region, and the resulted free electrons and holes will move towards opposite directions, forming the photocurrent in external circuit. When the light is turned off, the photo-generated electron-hole pairs disappear and the current decreases again. Since the photocurrent is governed by the graphene/p-type ZnTe:Ga NW junction, the surface effect will not significantly affect the photocurrent anymore. As a result, a large photoresponse is observed for the ZnTe:Ga NW/graphene Schottky diodes at reverse bias.

#### Conclusions

In summary, we present a synthetic method to prepare p-type ZnTe:Ga NWs by using a simple thermal evaporation method, in which Ga/Ga<sub>2</sub>O<sub>3</sub> powder was used as the dopant. Electrical analysis reveals that the Ga doping led to p-type conduction characteristics, and the conductivity could be tuned over four orders of magnitude by varying the doping level. It is also found that when the Ga content in the ZnTe NWs increases from 1.3 to 5.1–8.7%, the hole mobility and hole concentration will increase from 0.0069, to 0.33 and 0.46 cm<sup>2</sup> V<sup>-1</sup> s<sup>-1</sup> and from 2.4 × 10<sup>17</sup>, to

 $1.3 \times 10^{18}$  and  $5.6 \times 10^{18}$  cm<sup>-3</sup>, respectively. Moreover, optoelectronic characterization of photodetector composed of ZnTe:Ga NW/graphene Schottky diodes shows that the device exhibits high sensitivity to light illumination with good reproducibility at both forward and revers bias voltages. It is believed that the as-doped p-type ZnTe:Ga NWs with tunable electrical properties will have potential application for fabricating newgeneration nanoelectronic and nano-optoelectronic devices.

#### Acknowledgements

This work was supported by the Natural Science Foundation of China (NSFC, nos 21101051, 61106010, 21205026, 21401036), the Natural Science Foundation of Anhui Province of China (no. 140808MB31) and the Fundamental Research Funds for the Central Universities (2012HGCX0003, 2013HGCH0012, 2014HGCH0005).

#### References

- 1 J. S. Jie, W. J. Zhang, I. Bello, C. S. Lee and S. T. Lee, *Nano Today*, 2010, **5**, 313.
- 2 B. Nie, L. B. Luo, J. J. Chen, J. G. Hu, C. Y. Wu, L. Wang, Y. Q. Yu, Z. F. Zhu and J. S. Jie, *Nanotechnolgy*, 2013, 24, 095603.
- 3 C. Xie, L. B. Luo, L. H. Zeng, L. Zhu, J. J. Chen, B. Nie, J. G. Hu, Q. Li, C. Y. Wu and J. S. Jie, *CrystEngComm*, 2012, 14, 7222.
- 4 R. M. Ma, L. Dai, H. B. Huo, W. J. Xu and G. G. Qin, *Nano Lett.*, 2007, 7, 3300.
- 5 B. Nie, J. G. Hu, L. B. Luo, C. Xie, L. H. Zeng, P. Lv, F. Z. Li, J. S. Jie, M. Feng, C. Y. Wu, Y. Q. Yu and S. H. Yu, *Small*, 2013, 9, 2872.

- 6 Y. Dai, B. Yu, Y. Ye, P. C. Wu, H. Meng, L. Dai and G. G. Qin, J. Mater. Chem., 2012, 22, 18442.
- 7 F. Z. Li, L. B. Luo, Q. D. Yang, D. Wu, C. Xie, B. Xie, J. S. Jie,
   C. Y. Wu and S. H. Yu, *Adv. Energy Mater.*, 2013, 3, 579.
- 8 C. Soci, A. Zhang, B. Xiang, S. A. Dayeh, D. P. R. Aplin, J. Park,
   X. Y. Bao, Y. H. Lo and D. Wang, *Nano Lett.*, 2007, 7, 1003.
- 9 X. X. Wang, W. S. Zhang, G. L. Song, B. S. Zou, Z. S. Li, S. Guo, J. Li, Q. S. Li and R. B. Liu, *RSC Adv.*, 2014, 4, 52550.
- 10 L. B. Luo, X. L. Huang, M. Z. Wang, C. Xie, C. Y. Wu, J. G. Hu, L. Wang and J. A. Huang, *Small*, 2014, **13**, 2645.
- 11 M. C. Schlamp, X. G. Peng and A. P. Alivisatosa, *J. Appl. Phys.*, 1997, **82**, 5837.
- 12 J. I. Wong, N. Mishra, G. C. Xing, M. J. Li, S. Chakrabortty, T. C. Sum, Y. M. Shi, Y. T. Chan and H. Y. Yang, *ACS Nano*, 2014, 8, 2873.
- 13 S. Lee, H. Kim, D. J. Yun, S. W. Rhee and K. Yong, *Appl. Phys. Lett.*, 2009, **95**, 262113.
- 14 C. Xie, B. Nie, L. Zhu, L. H. Zeng, Y. Q. Yu, X. H. Wang, Q. L. Fang, L. B. Luo and Y. C. Wu, *Nanotechnology*, 2013, 24, 355203.
- 15 J. M. Wu and C. C. Kao, RSC Adv., 2014, 4, 13882.
- 16 L. H. Zhang, L. L. Fan, Z. Li, E. Z. Shi, X. M. Li, H. B. Li, C. Y. Ji, Y. Jia, J. Q. Wei, K. L. Wang, H. W. Zhu, D. H. Zhu and A. Y. Cao, *Nano Res.*, 2011, 4, 891.
- 17 A. Goudarzi, A. D. Namghi and C. S. Ha, *RSC Adv.*, 2014, 4, 59764.
- 18 A. Erlacher, M. Ambrico, G. Perna, L. Schiavulli, T. Ligonzo, H. Ligonzo and B. Ullrich, *Appl. Surf. Sci.*, 2005, 248, 402.
- 19 Q. Zhang, J. Zhang, M. I. B. Utama, B. Peng, M. de la Mata, J. Arbiol and Q. H. Xiong, *Phys. Rev. B: Condens. Matter Mater. Phys.*, 2012, 85, 085418.
- 20 E. Groeneveld, S. V. Berkum, A. Meijerink and C. D. M. Donegá, *Small*, 2011, 7, 1247.
- 21 K. P. Acharya, A. Erlacher and B. Ullrich, *Thin Solid Films*, 2007, **515**, 4066.
- 22 A. E. Rakhshani and S. Thomas, J. Mater. Sci., 2013, 48, 6386.

- 23 Q. Wu, M. Litz and X. C. Zhang, *Appl. Phys. Lett.*, 1996, **68**, 2924.
- 24 T. Tanaka, K. M. Yu, P. R. Stone, J. W. Beeman, O. D. Dubon, L. A. Reichertz, V. M. Kao, M. Nishio and W. Walukiewicz, *J. Appl. Phys.*, 2010, **108**, 024502.
- 25 L. Li, Y. W. Yang, X. H. Huang, G. H. Li and L. D. Zhang, J. Phys. Chem. B, 2005, **109**, 12394.
- 26 E. Janik, P. Dluzewski, S. Kret, A. Presz, H. Kirmse,
  W. Neumann, W. Zaleszczyk, L. T. Baczewski,
  A. Petroutchik, E. Dynowska, J. Sadowski, W. Caliebe,
  G. Caliebe and T. Wojtowicz, *Nanotechnology*, 2007, 18, 475606.
- 27 D. Wu, Y. Jiang, Y. G. Zhang, Y. Q. Yu, Z. F. Zhu, X. Z. Lan, F. Z. Li, C. Y. Wu, L. Wang and L. B. Luo, *J. Mater. Chem.*, 2012, 22, 23272.
- 28 Z. Li, J. Salfi, C. D. Souza, P. Sun, S. V. Nair and H. E. Ruda, *Appl. Phys. Lett.*, 2010, **97**, 063510.
- 29 L. Li, Y. W. Yang, X. H. Huang, G. H. Li and L. D. Zhang, J. Phys. Chem. B, 2005, **109**, 12394.
- 30 S. Y. Li, Y. Jiang, D. Wu, L. Wang, H. H. Zhong, B. Wu, X. Z. Lan, Y. Q. Yu, Z. B. Wang and J. S. Jie, *J. Phys. Chem. C*, 2010, **114**, 7980.
- 31 H. B. Huo, L. Dai, C. Liu, L. P. You, W. Q. Yang, R. M. Ma, G. Z. Ran and G. G. Qin, *Nanotechnology*, 2006, 17, 5912.
- 32 M. Z. Wang, W. J. Xie, H. Hu, Y. Q. Yu, C. Y. Wu, L. Wang and L. B. Luo, *Appl. Phys. Lett.*, 2013, **103**, 213111.
- 33 P. A. Hu, Z. Z. Wen, L. F. Wang, P. H. Tan and K. Xiao, *ACS Nano*, 2014, **6**, 5988.
- 34 B. Y. Zu, B. Lu, Y. N. Guo and X. C. Dou, *J. Mater. Chem. C*, 2014, **2**, 2045.
- 35 Y. Q. Yu, J. S. Jie, P. Jiang, L. Wang, C. Y. Wu, Q. Peng, X. W. Zhang, Z. Wang, C. Xie, D. Wu and Y. Jiang, *J. Mater. Chem.*, 2011, 21, 12632.
- 36 D. Wu, Y. Jiang, X. D. Yao, Y. J. Chang, Y. G. Zhang, Y. Q. Yu,
  Z. F. Zhu, Y. Zhang, X. Z. Lan and H. H. Zhong, *J. Mater. Chem. C*, 2014, 2, 6547.